

# Coherence, Correlation, convolution and matched-filtering

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What keeps radio astronomers busy...  
OR what are they after really...mostly.?

Sky Noise !, which brings us the cosmic news,  
unlike in communications, where signal and noise  
have explicit distinction.

Noise properties (necessarily statistical)....  
related to color (spectrum), prob. density fn.  
coherence, if any (spatial, temporal)  
and their variation in time, and polarization state,  
etc.

What keeps radio astronomers busy...  
OR what are they after really...mostly.?

Intensity (i.e. variance of noise)(1) as a function of

- Sky co-ordinates (2,3)
- Frequency (4)
- Time (5)
- Polarization state (6)

The spans and resolutions in these six parameters,  
i.e. the 12 numbers, together would essentially  
define vital-statistics of a telescope.

But not all can be specified mutually  
independently !

# But not all can be specified mutually independently !

Intensity : span/resolution : dynamic range

: dependence on collecting area,

spectral bandwidth, observation time (duration)

sky coverage, angular resolution:

depend on collecting area & frequency

Frequency : resolution limited to total time-span

Time : resolution limited to total spectral-span

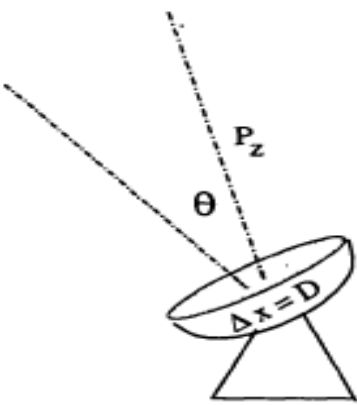
Polarization state : antenna response in polarization

# Transducer for sensing incident radiation at radio frequencies

- Individual photon energies ( $h\nu$ ) are too small at radio frequencies, hence even a swarm of millions of radio photons can not free an electron from its atom, or initiate photo-chemical changes.
- The radio photons do however induce currents in antenna. The currents are too weak to trigger any detector, unless can be amplified in suitable device (e.g. by amplifiers).

# Unavoidables and relevant in even the single photon case

**A** Take a telescope of diameter  $D$  operating at any wavelength  $\lambda$   
Let its beamwidth be  $\theta$



Now uncertainty of photon arrival position is clearly  $\Delta x = D$

Uncertainty of photon transverse momentum  $\Delta P_x$  is  $P_z \theta$

$$\Delta x \Delta P_x = D P_z \theta \sim h, \text{ or } \theta \approx \frac{h}{D P_z}$$

now  $P_z = \frac{\text{energy}}{\text{velocity}} = \frac{h\nu}{c}$

$$\theta \approx \frac{h c}{h \nu D} = \frac{\lambda}{D}$$

**YOUR DIFFRACTION FORMULA !!**

**B** Similarly,  $\Delta z \Delta P_z \sim h$

$\Delta z =$  Uncertainty in longitudinal position  
 $= c \Delta t$

And  $\Delta P_z =$  Uncertainty in longitudinal momentum

But longitudinal momentum  $= \frac{h\nu}{c}$

so  $\Delta P_z = \Delta\nu \frac{h}{c}$

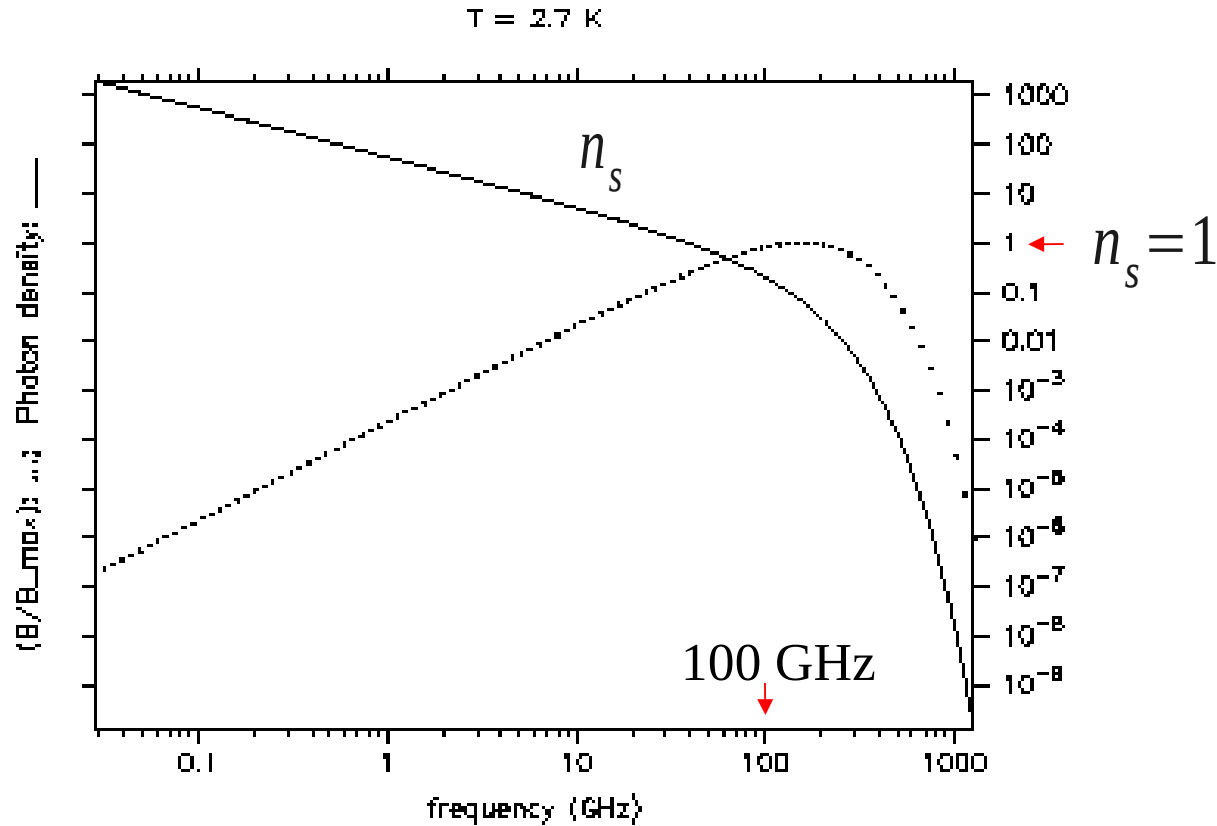
so,  $\Delta z \Delta P_z \sim h \longrightarrow c \Delta t \Delta\nu \frac{h}{c} \sim h$

or  $\Delta\nu \approx \frac{1}{\Delta t}$

i.e. **RESPONSE TIME OF A FILTER !!**

**Figure 33-3.** Heisenberg's uncertainty principle connects naturally, A) the aperture size and resolving power of telescopes, or B) the bandwidths and response times of filters, even when dealing with single photons.

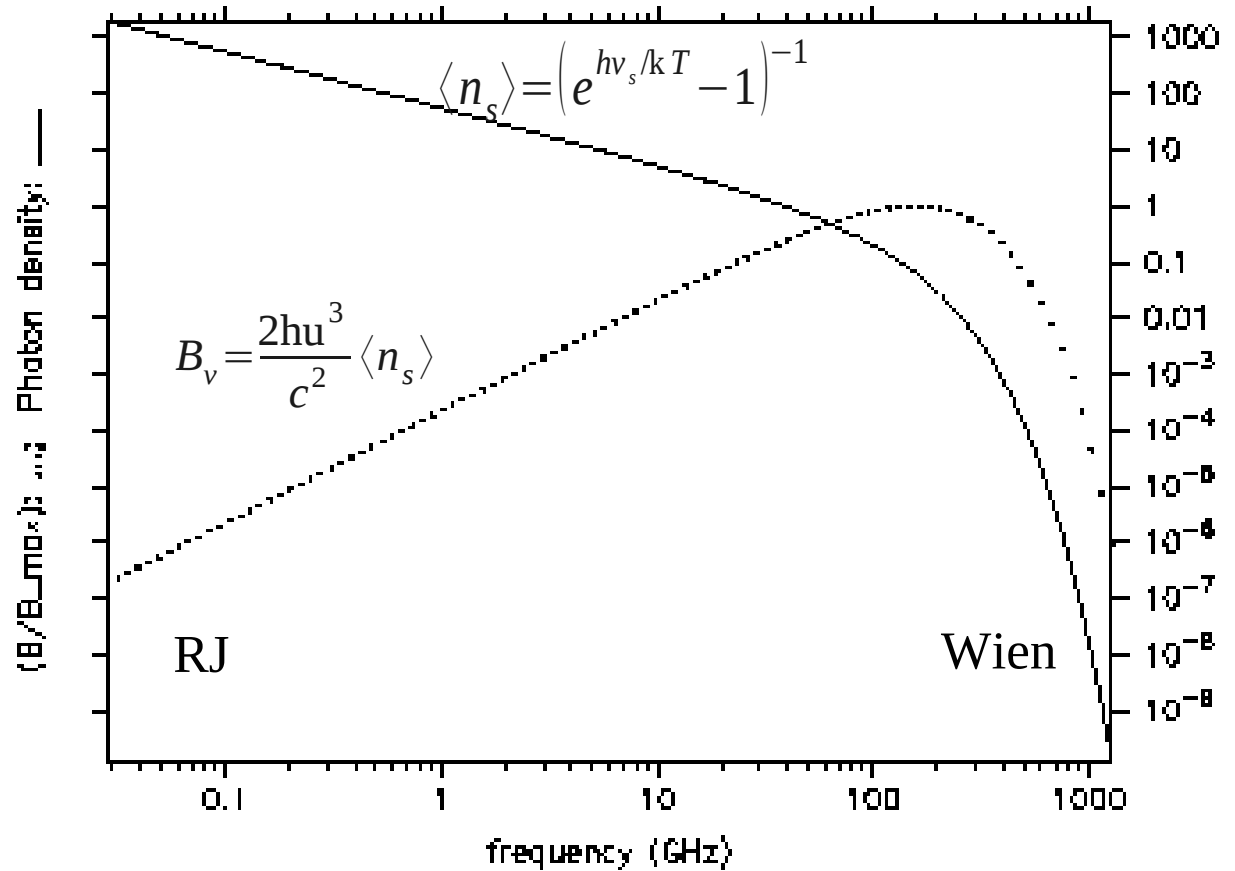
# The sky is not dark in the radio!



“Even the feeble microwave background ensures that the occupation number at most radio frequencies is already high. In other words, even though the particular contribution to the signal that we seek is very very weak, it is already in a **classical** sea of noise and if there are benefits to be derived from retaining the associated aspects, we would be foolish to pass them up.” Radhakrishnan 1998

T = 2.7 K

When is wave noise important? Photon occupation number at 2.7K



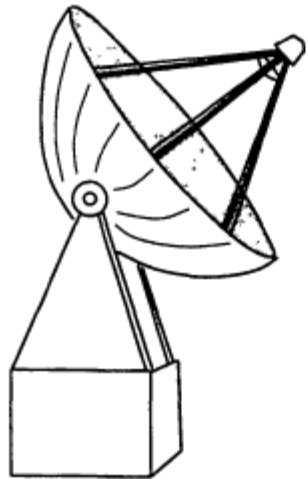
$$\langle Dn_s^2 \rangle^0 \langle (n_s - \langle n_s \rangle)^2 \rangle = \langle n_s \rangle + \langle n_s \rangle^2$$

Wien:  $n_s < 1 \Rightarrow \text{rms } \mu \sqrt{n_s}$  (countingstats)

RJ:  $n_s > 1 \Rightarrow \text{rms } \mu n_s$  (wavenoise)



A RADIO



IF source is strong,  $T_A \gg T_{REC}$

$$T_A \propto D^2$$

$$\Delta T_A = \frac{T_A}{\sqrt{t\Delta\nu}}$$

$$\therefore \frac{\Delta T_A}{T_A} = \frac{1}{\sqrt{t\Delta\nu}}$$

INDEPENDENT of DIAMETER.

B OPTICAL



$N$  (no. of photons)  $\propto D^2$

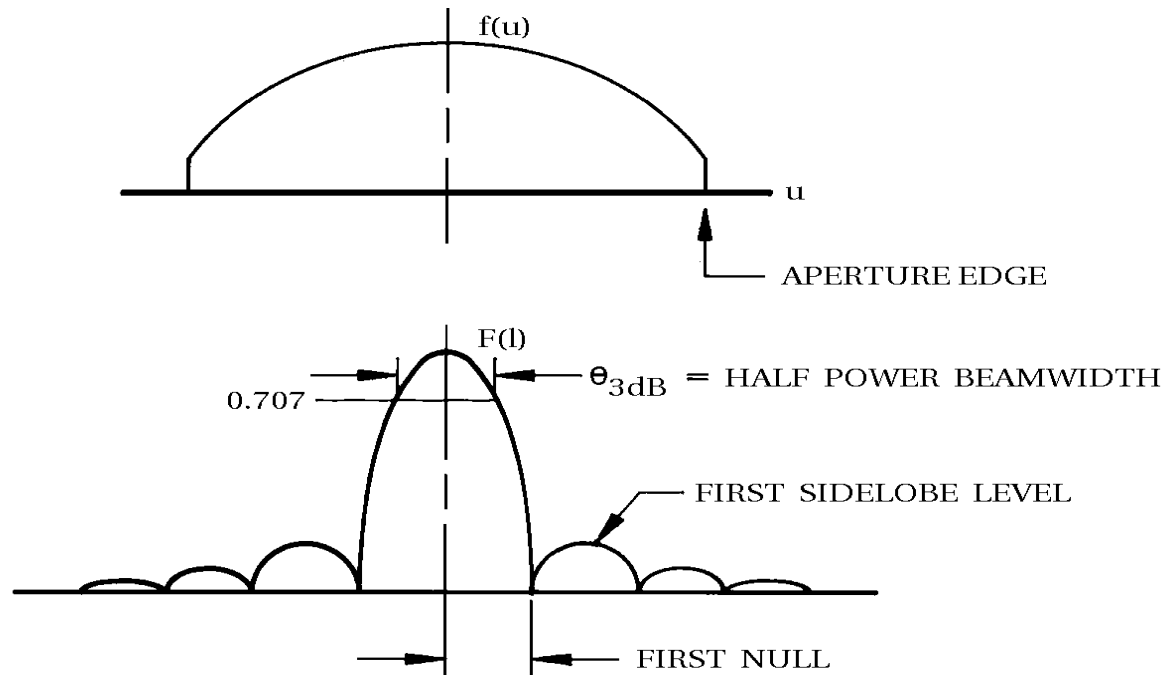
$$\frac{\Delta N}{N} = \frac{1}{\sqrt{N}}$$

$$\therefore \frac{\Delta N}{N} \propto \frac{1}{D}$$

THE BIGGER THE BETTER, ALWAYS.

# Fourier pairs in astronomy:

aperture illumination  $\leftrightarrow$  far-field diffraction pattern



Another Fourier pairs in astronomy:

brightness distribution  $\leftrightarrow$  spatial-coherence  
function

# Fourier pairs in astronomy:

aperture illumination  $\leftrightarrow$  far-field diffraction pattern

Recall FT stuff

Shift  $\leftrightarrow$  phase gradient

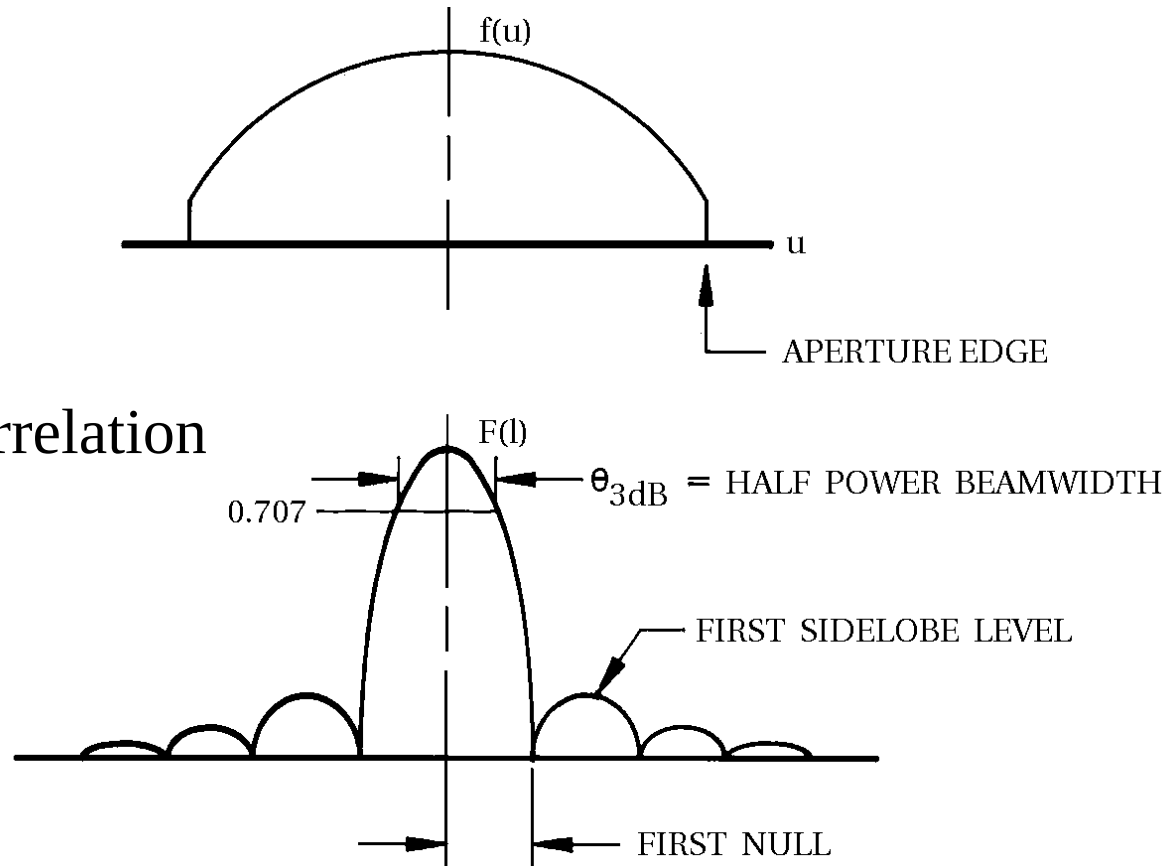
Product  $\leftrightarrow$  convolution

Mod square  $\leftrightarrow$  auto-correlation

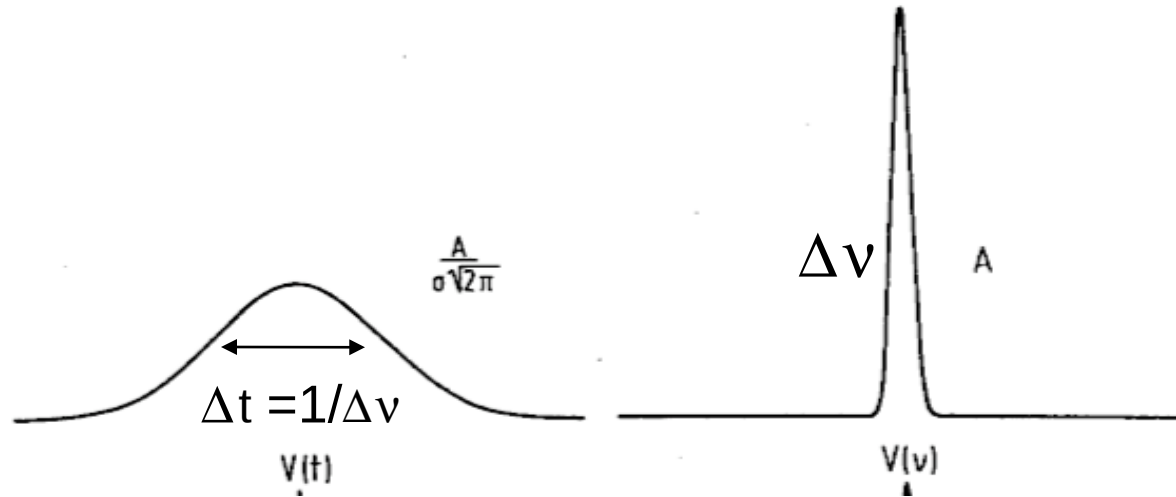
Also, linear combinations  
preserved across the  
two domains

Same is relevant for the  
other Fourier pair

Time-sequence & spectrum



# General Fourier conjugate variable relationships

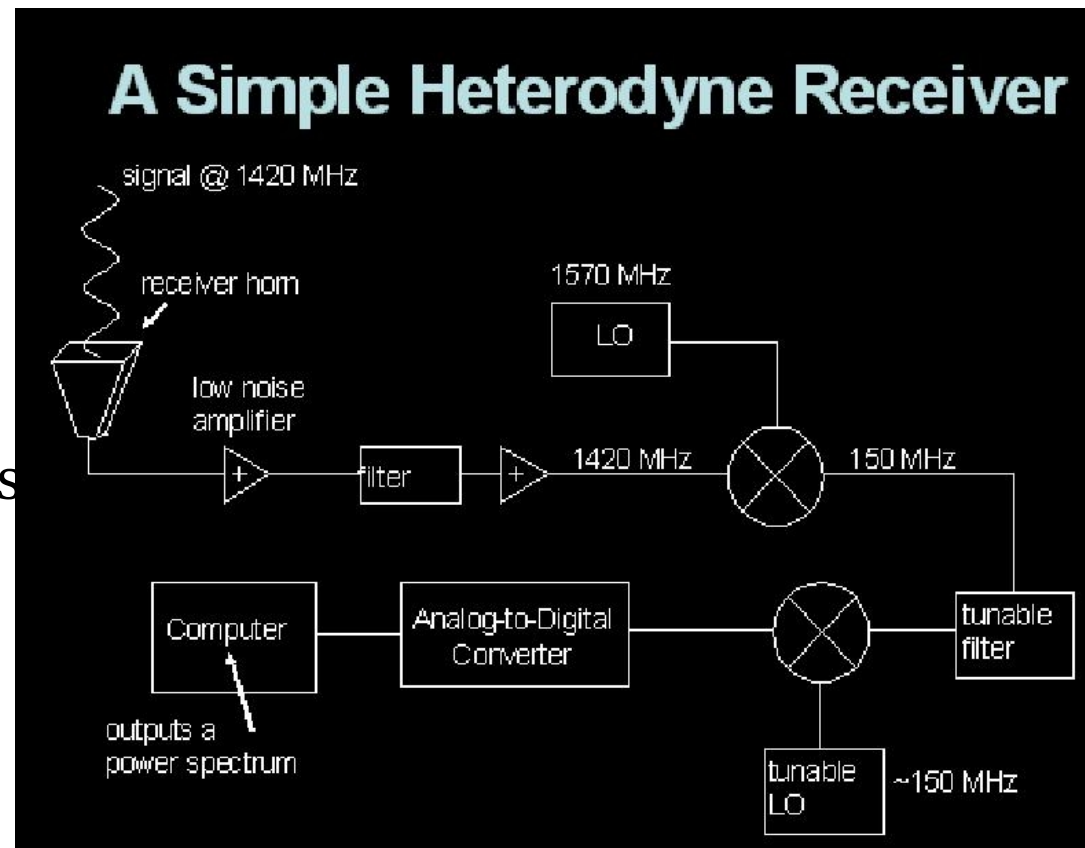


- Fourier conjugate variables, frequency -- time (or power spectrum in freq, autocorrelation in lag, eg. Wiener-Khinchin theorem)
- If  $V(v)$  is Gaussian of width  $\Delta v$ , then  $V(t)$  is also Gaussian of width =  $\Delta t = 1/\Delta v$
- **Measurements of  $V(t)$  on timescales  $\Delta t < 1/\Delta v$  are correlated, ie. *not independent***
- **Restatement of Nyquist sampling theorem: maximum information is gained by sampling at  $\sim 1/2\Delta v$ . Nothing changes on shorter timescales.**

# To be revisited later, but mention key aspects: noise from receiver (Trec) and “coherent” amplification

Desired : amplification  
Price paid: additional  
noise from  
the amplifier  
and other  
receiver components

even though one tries to keep  
these costs to minimum  
e.g. low-noise amplifiers, etc.



# Noise limit: quantum noise and coherent amplifiers

Uncertainty principle for photons:

$$DE \, Dt = h$$

$$DE = h\nu Dn_s$$

$$Dt = \frac{Dj}{\nu 2p}$$

$$\Rightarrow Dj Dn_s = 1 \text{ rad Hz}^{-1} \text{sec}^{-1}$$

Coherent Amplifier:  $Dj < 1 \text{ rad} \Rightarrow Dn_s = 1 \text{ photon Hz}^{-1} \text{sec}^{-1}$

Phase coherent amplifier has minimum noise of  $n_s = 1 \text{ photon Hz}^{-1} \text{sec}^{-1}$

**Phase coherent amplifier automatically puts signal into RJ regime => wave noise dominated**

Note: phase coherent amplifier is not a detector

