



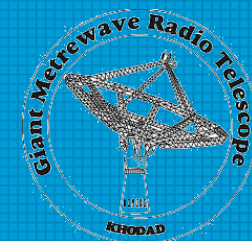
# Time-domain (Pulsar/FRB) Observing Techniques

**Yogesh Maan**

## Astronomical Techniques II

Main references and sources:

- Low-frequency Radio Astronomy (primarily Chapters 6 and 17),
- Handbook of Pulsar Astronomy,
- J. van Leeuwen's talks available online; and internet.



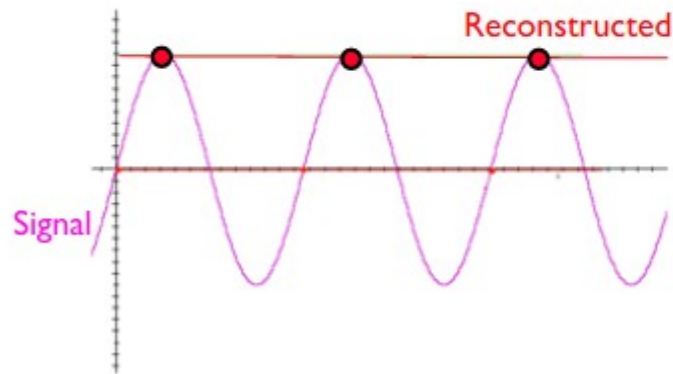
# Sensitivity: The Radiometer Equation

- The minimum temperature that a telescope can measure is limited by the noise (root mean square) fluctuations in the receiver system.

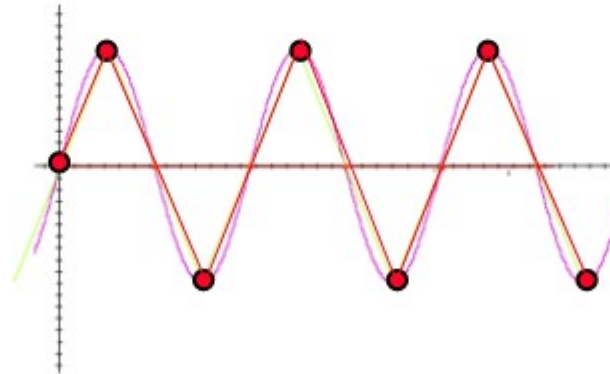
$$\Delta T_{\text{sys}} = \frac{T_{\text{sys}}}{\sqrt{n_p t \Delta f}}$$

# Sensitivity: The Radiometer Equation

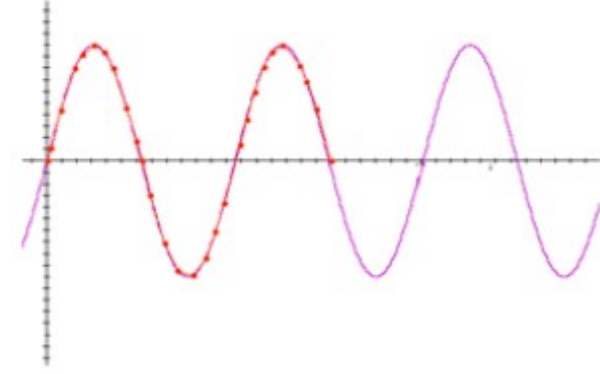
If we sample once per cycle time (period) we would consider the signal to have a constant amplitude.



If we sample twice per cycle time (period) we get a saw-tooth wave that is becoming a good approximation to a sinusoid.



For lossless digitisation we must sample the signal at least twice per cycle time.



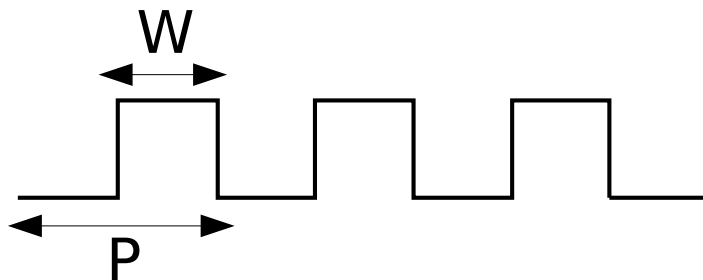
- Nyquist sampling theorem: Number of samples per second  $> 2xf_{\max}$  (2xBandwidth)

# Sensitivity: The Radiometer Equation

- The minimum detectable “mean” flux density (corresponding to a S/N threshold) depends on receiver noise as well as pulse-width relative to the pulsar’s rotation period.

$$S = \frac{2k_{\text{B}}T_{\text{A}}}{A_{\text{e}}} = \frac{T_{\text{A}}}{G}$$

$$S_{\text{min}} = \beta \frac{(S/N_{\text{min}})T_{\text{sys}}}{G \sqrt{n_{\text{p}} t_{\text{int}} \Delta f}} \sqrt{\frac{W}{P - W}}$$



(Handbook of Pulsar Astronomy, A1.4)

# Sensitivity: The Radiometer Equation

$$S = \frac{2k_B T_A}{A_e} = \frac{T_A}{G} \quad S_{\min} = \beta \frac{(S/N_{\min}) T_{\text{sys}}}{G \sqrt{n_p t_{\text{int}} \Delta f}} \sqrt{\frac{W}{P - W}}$$



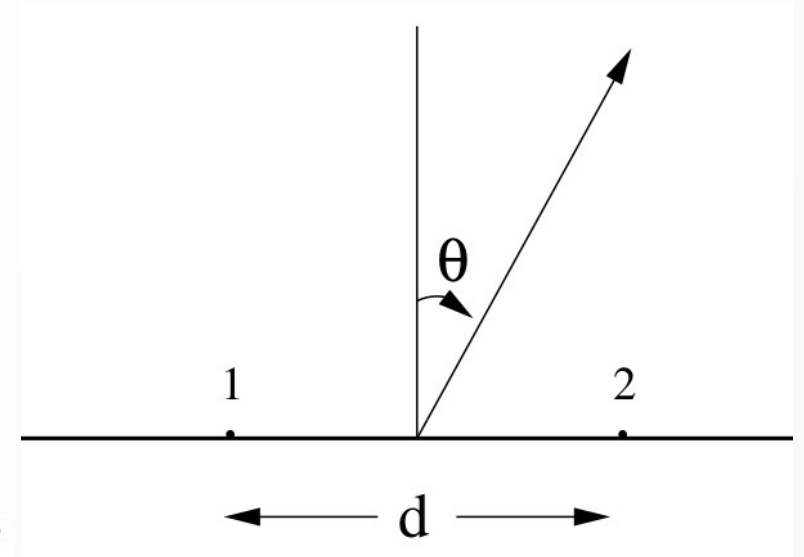
# Phased Arrays

- A two-elements phased array

- Reciprocity theorem: Performance of an antenna when collecting radiation from a point source at infinity may be studied by considering its properties as a *transmitter*:

$$E(\theta) = E_1 e^{j\psi/2} + E_2 e^{-j\psi/2}, \quad \psi = k d \sin \theta + \delta$$

$$E(\theta) = 2 E_0 \cos(\psi/2)$$



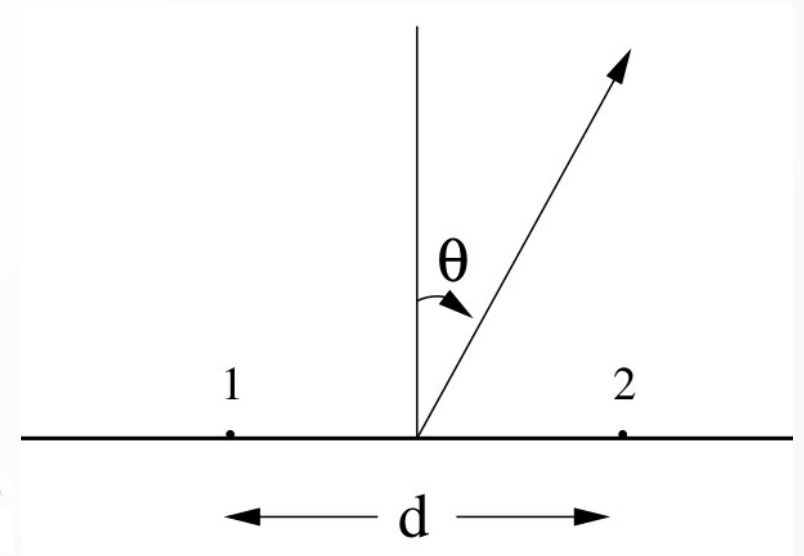
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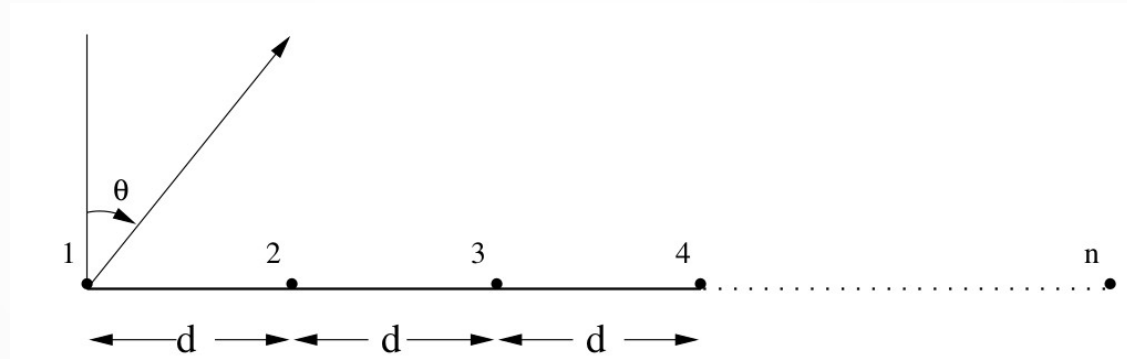
$$E(\theta) = 2 E_0 \cos(\psi/2)$$



- For  $d \gg \lambda$ , the field pattern is a sinusoidal function of  $\theta$ , with a period of  $2\lambda/d$ .
- $\delta \neq 0$  shifts the phase of the above pattern.
- The field pattern is weighted by the directional pattern of individual elements.

# Phased Arrays

- n-elements (equally spaced) phased array



$$E(\theta) = E_0 \left[ 1 + e^{j\psi} + e^{j2\psi} + \dots + e^{j(n-1)\psi} \right]$$

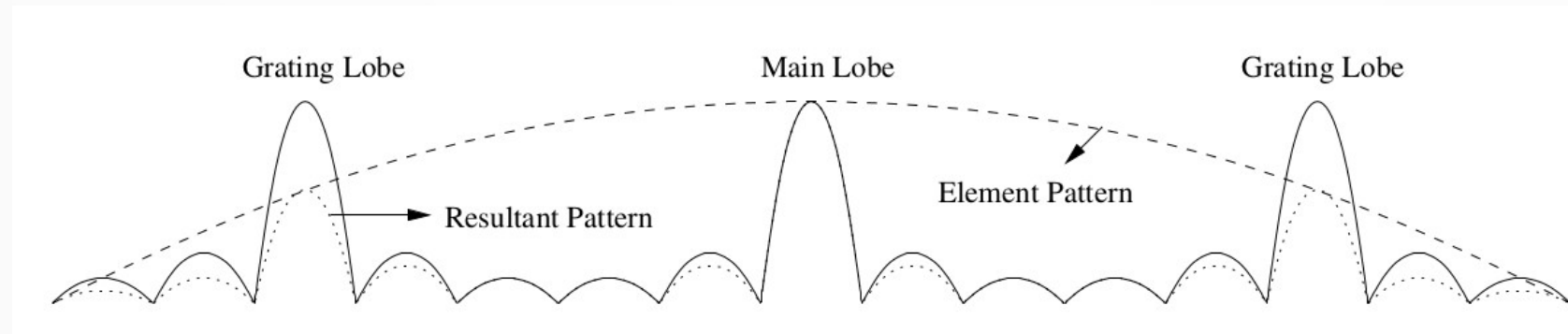
$$\psi = k d \sin \theta + \delta$$

$$E(\theta) = E_0 \frac{\sin(n\psi/2)}{\sin(\psi/2)} e^{j(n-1)\psi/2}$$



# Phased Arrays

- n-elements (equally spaced) phased array

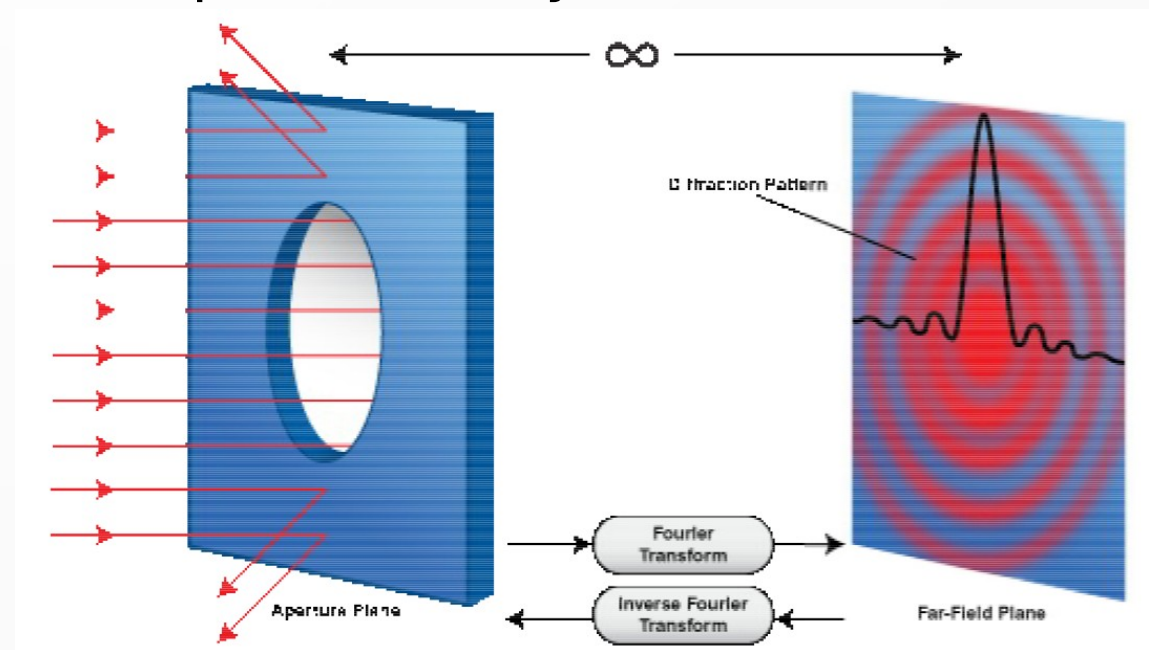


$$E(\theta) = E_0 \frac{\sin(n\psi/2)}{\sin(\psi/2)} e^{j(n-1)\psi/2} \quad \psi = kd \sin \theta + \delta$$

- For  $d > \lambda$  and  $\delta = 0$ , the field pattern is a periodic function of  $\theta$ , with maxima at  $\psi = 0, 2\pi, 4\pi, \dots$
- $\delta \neq 0$  shifts the phase of the above pattern.
- The field pattern is weighted by the directional pattern of individual elements.
- HPBW  $\approx \lambda/nd$ .

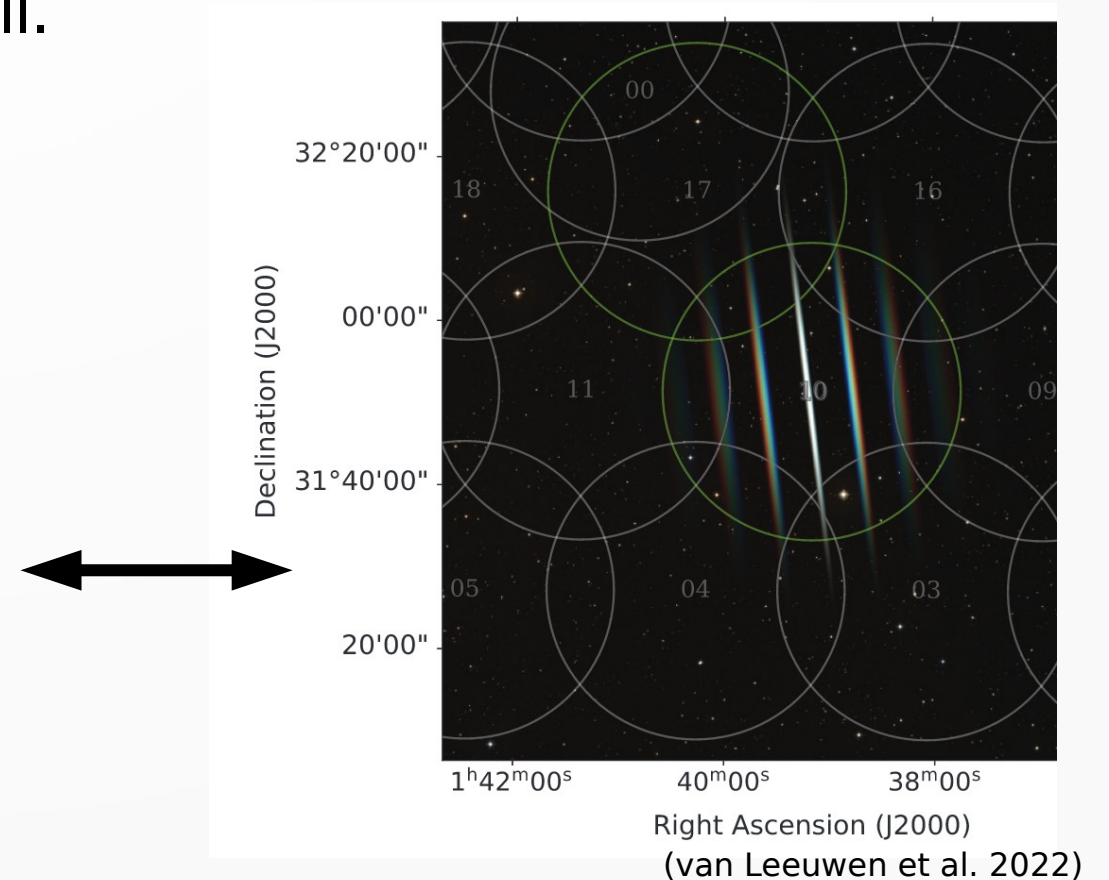
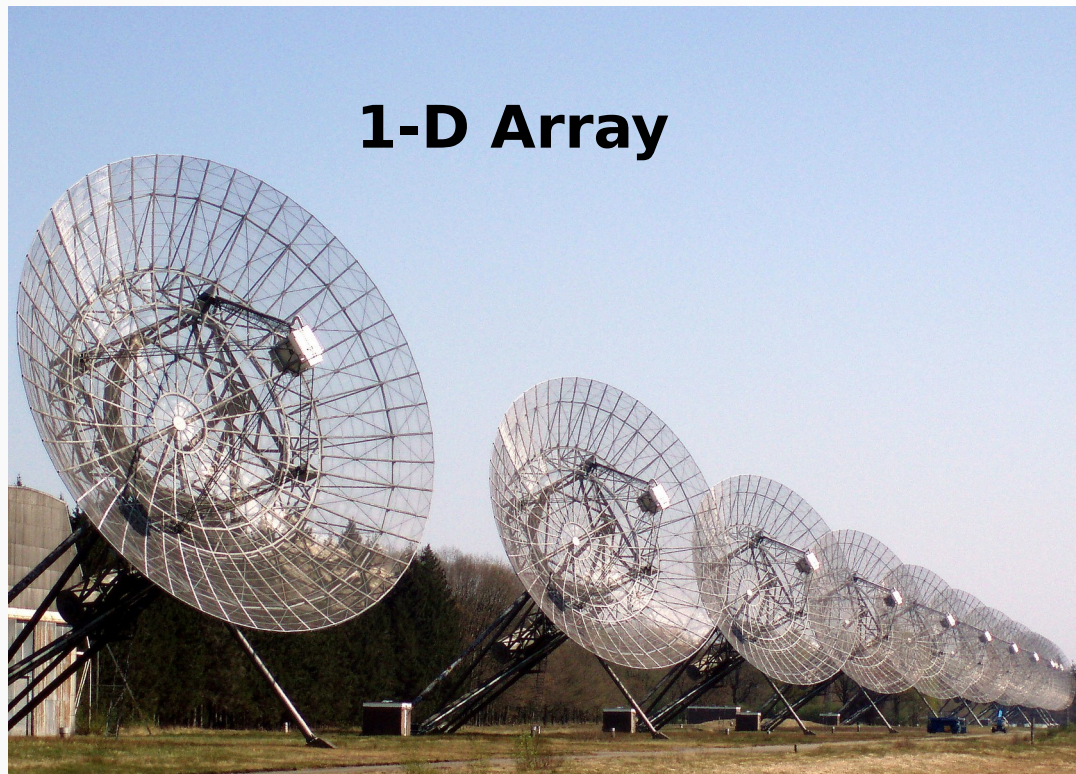
# Phased Arrays

- Two dimensional phased arrays
- Far-field radiation pattern (the antenna's "beam") is the Fourier transform of the aperture plane electric field distribution. True for "phased-array" radiation pattern as well.
- Radiation pattern of the 1-D n-element phased-array?



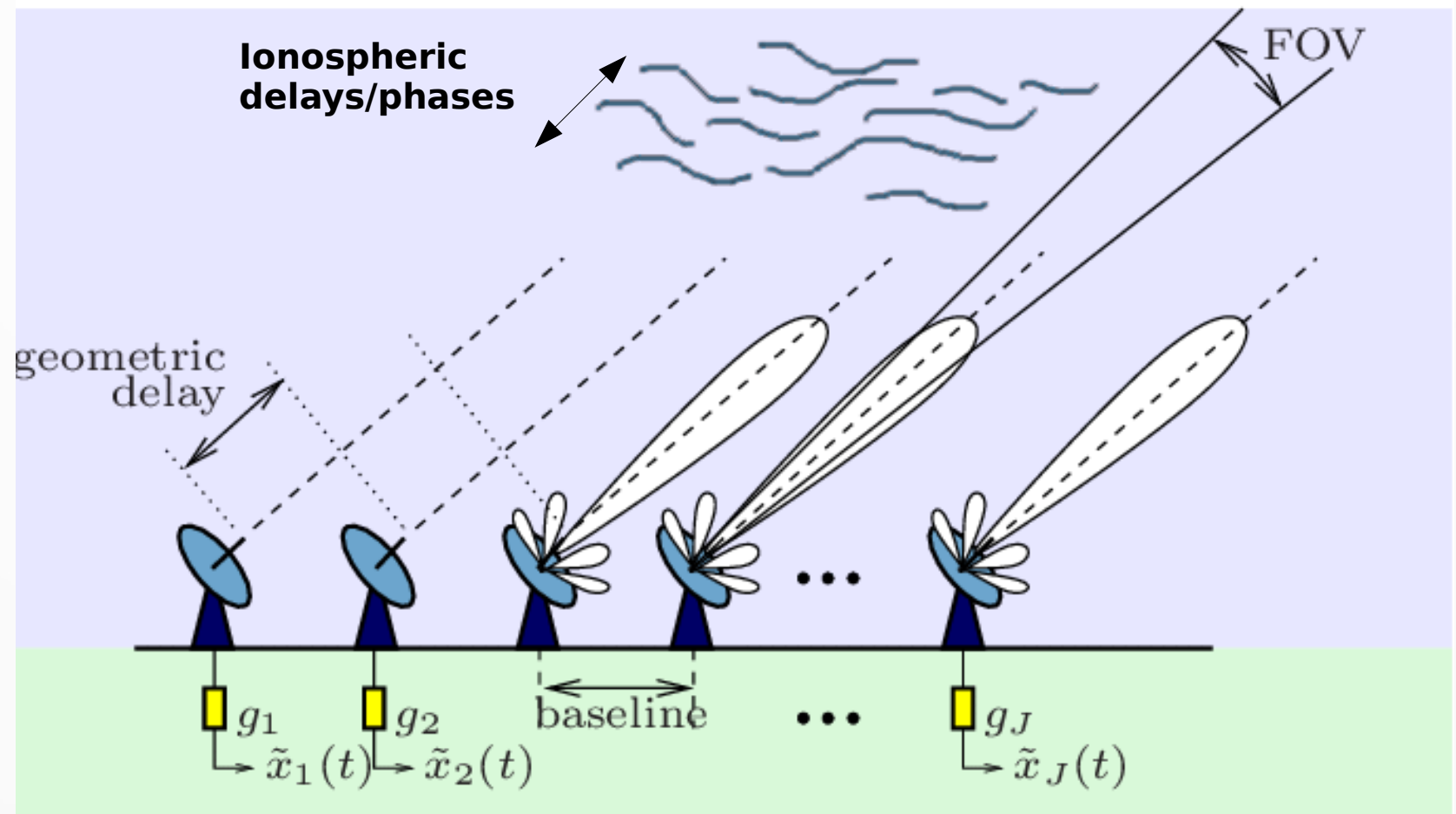
# Phased Arrays

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# Coherently Phased Arrays

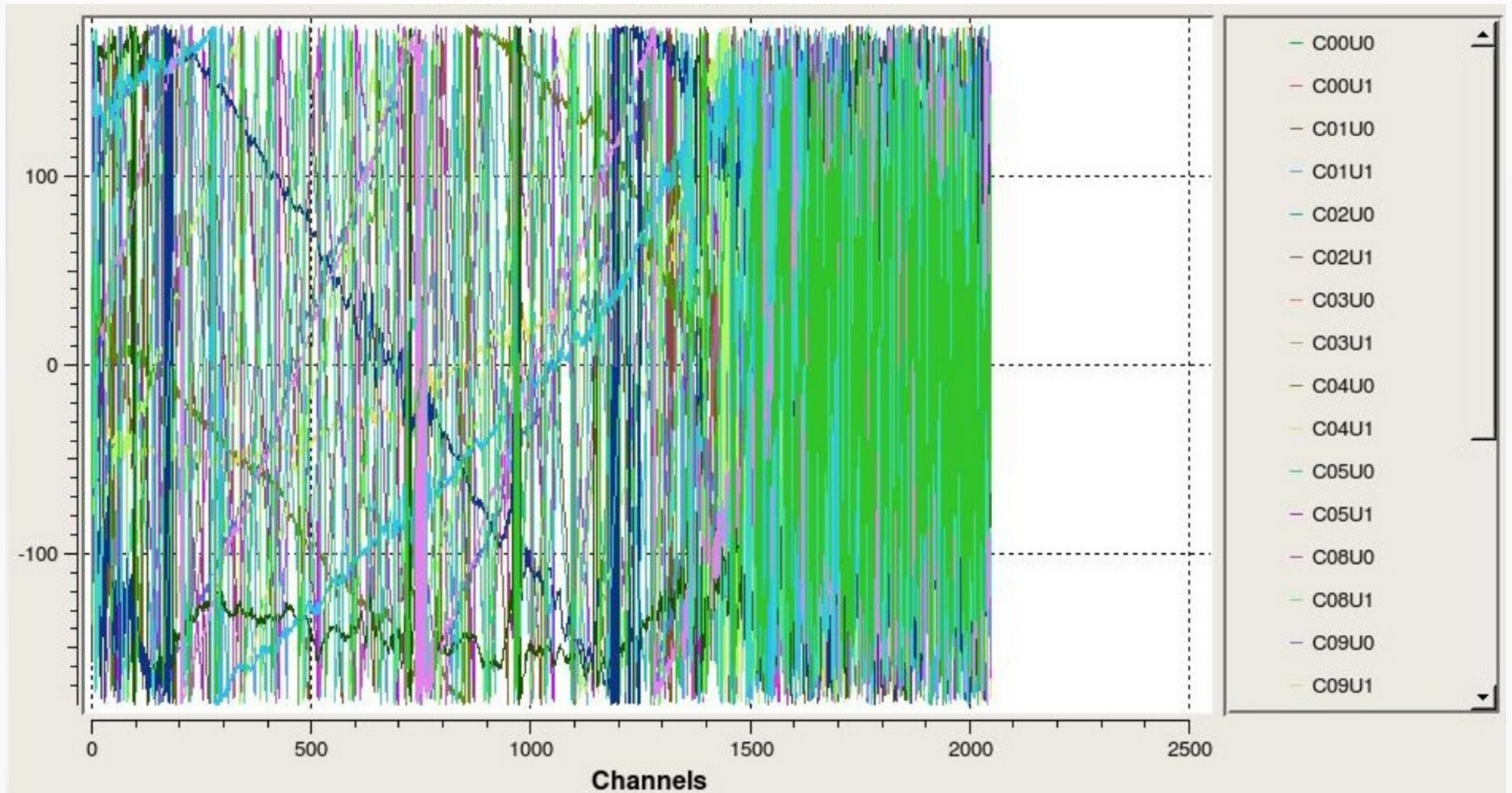
- Combine the **voltage** signals from the phased-array after all the delay corrections.
- $G \approx n \times G_0$
- Narrow beam



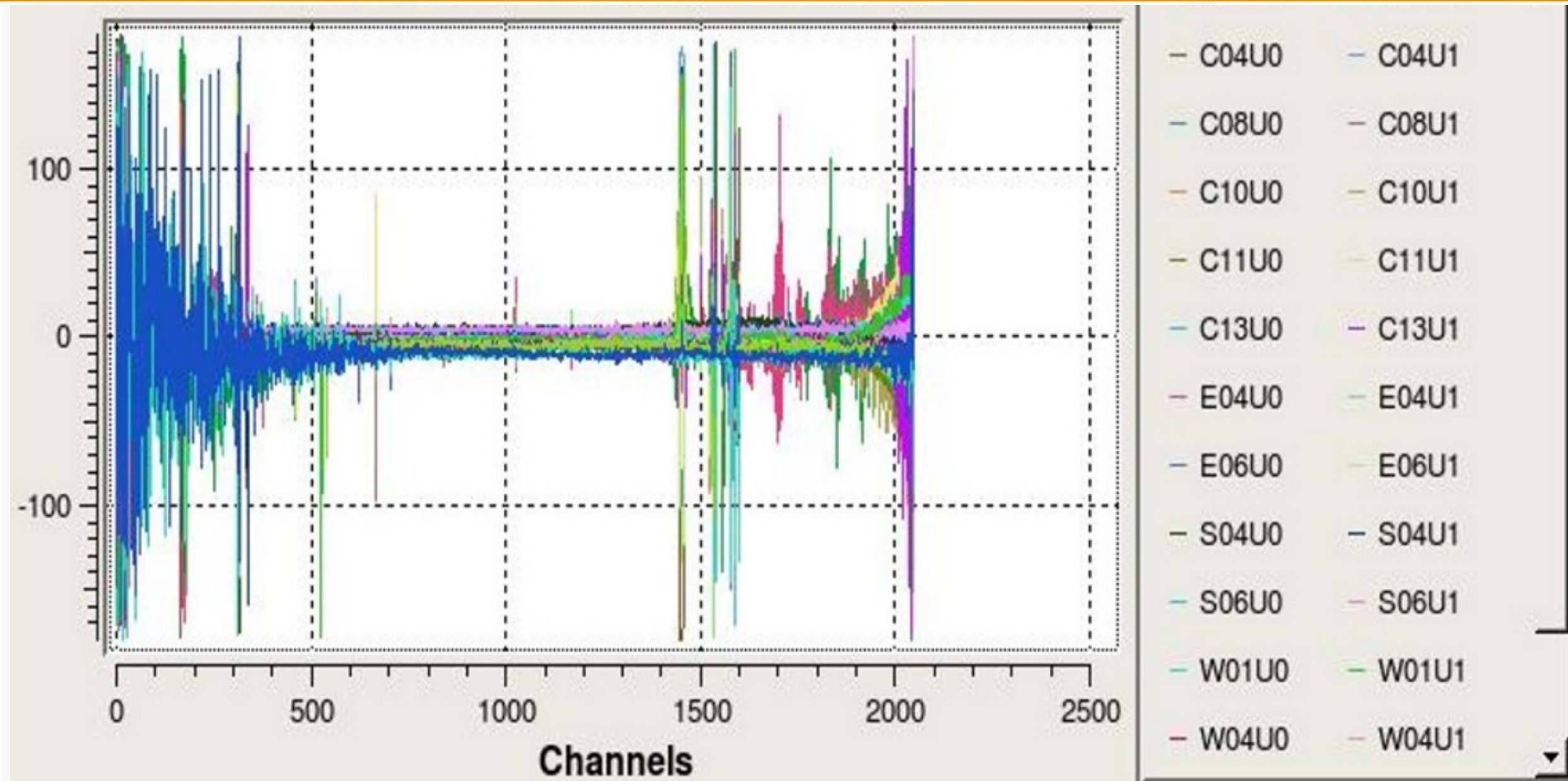
**Instrumental  
delays**

Figure adapted from "Signal processing tools for Radio Astronomy"

# Coherently Phased Arrays

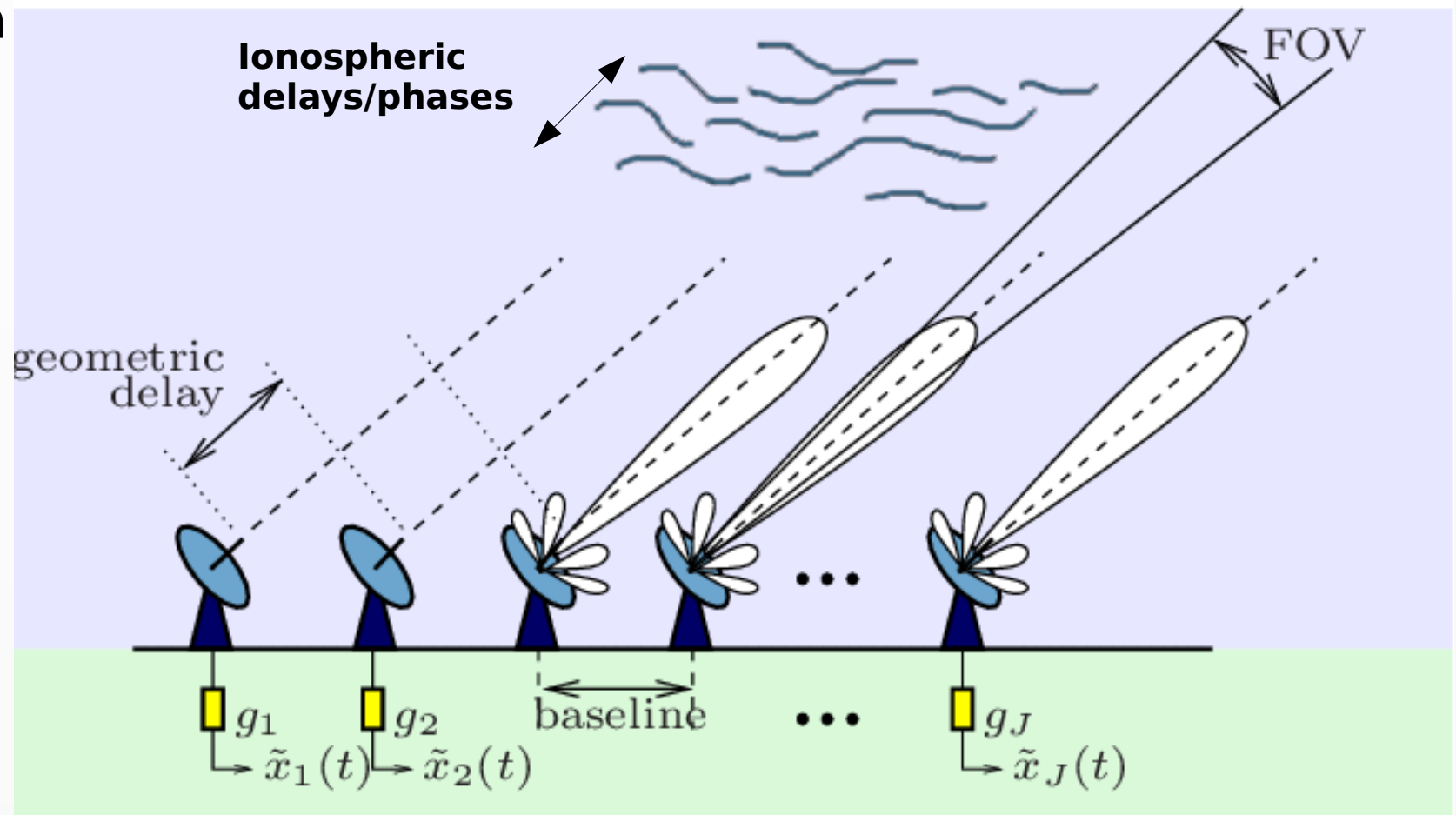


# Coherently Phased Arrays



# Incoherently Phased Arrays

- Combine the **power** signals from the phased-array after all the delay corrections.
- $G \approx \sqrt{n} \times G_0$
- Beam same as that of the primary element.



**Instrumental  
delays**

Figure adapted from "Signal processing tools for Radio Astronomy"

# Pulsar observations: Dedispersion

***Interstellar medium is (mostly) cold, ionized plasma***

$$\text{Refractive index: } \mu = [1 - (\nu_p/\nu)^2]^{1/2}$$

$$\text{Plasma frequency: } \nu_p = \left(\frac{e^2 n_e}{\pi m_e}\right)^{1/2}$$

$$\text{Group velocity: } v_g = \mu c$$

$$\text{Delay: } \Delta t = \int_0^d \frac{dl}{v_g} = \frac{e^2}{2\pi m_e c} \frac{\int_0^d n_e dl}{\nu^2}$$

$$\text{Dispersion Measure: } \text{DM} = \int_0^d n_e dl$$

***=> Frequency-dependent index of refraction implies different frequency signals propagate with different velocities.***

***=> Arrival time varies as a function of frequency***

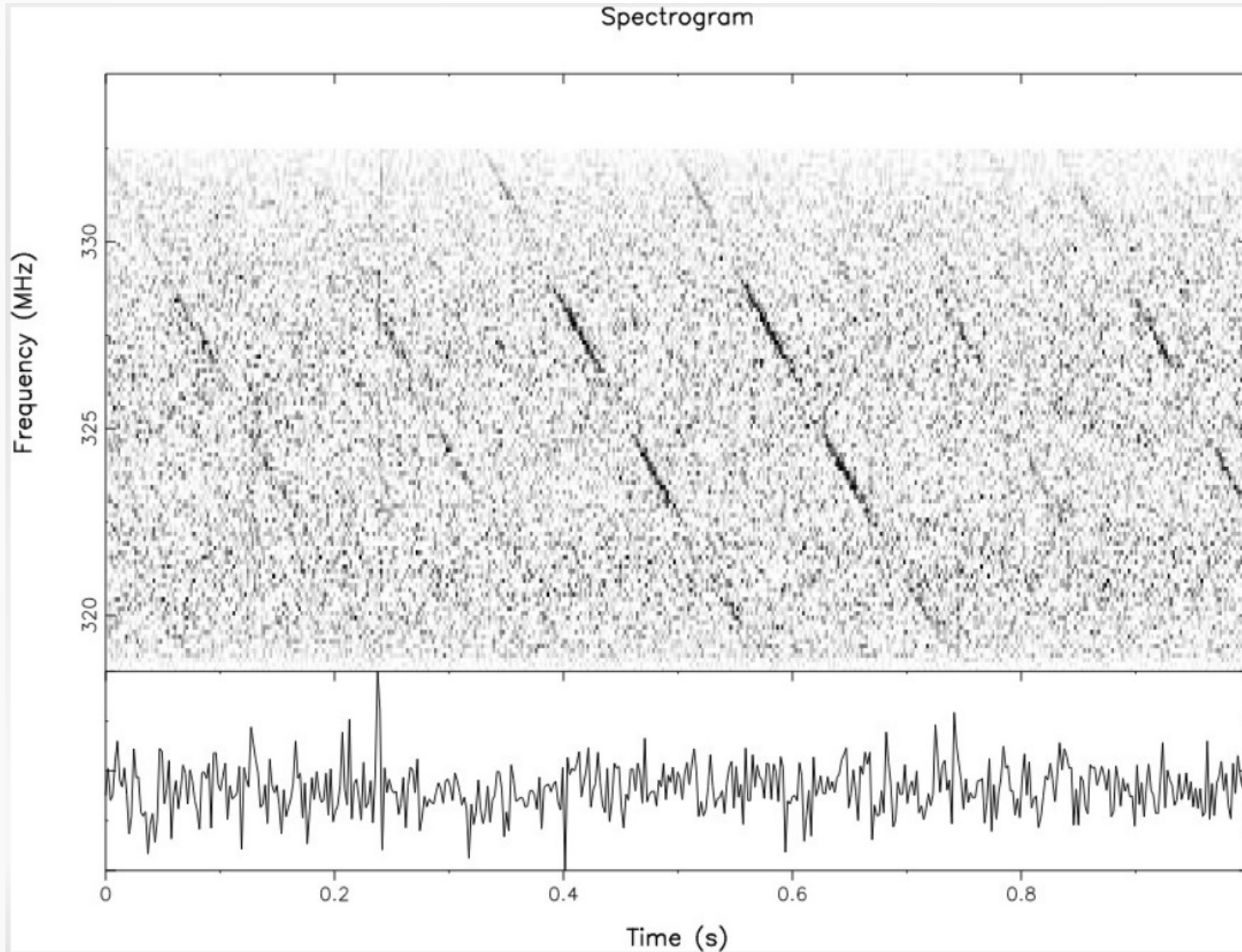


# Pulsar observations: Dedispersion

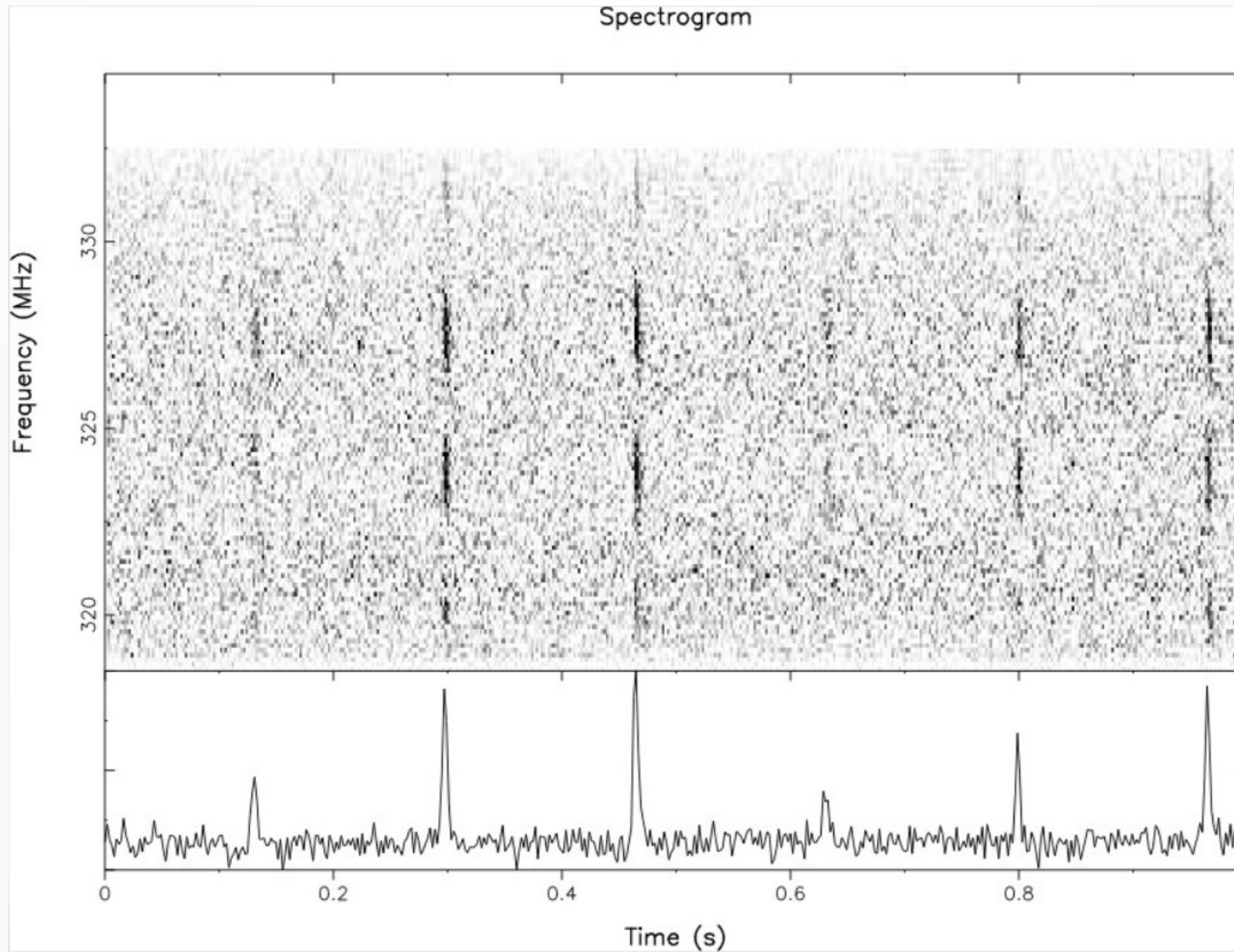
$$\Delta t(\text{ms}) = 4.15 \times 10^6 \times \int n_e d\ell \times (\nu_1^{-2} - \nu_2^{-2}); \quad \nu \text{ is in MHz}$$

$\int n_e d\ell \rightarrow$  **Dispersion Measure (DM)**  
(*Electron density integrated over the distance from the source to the observer*)

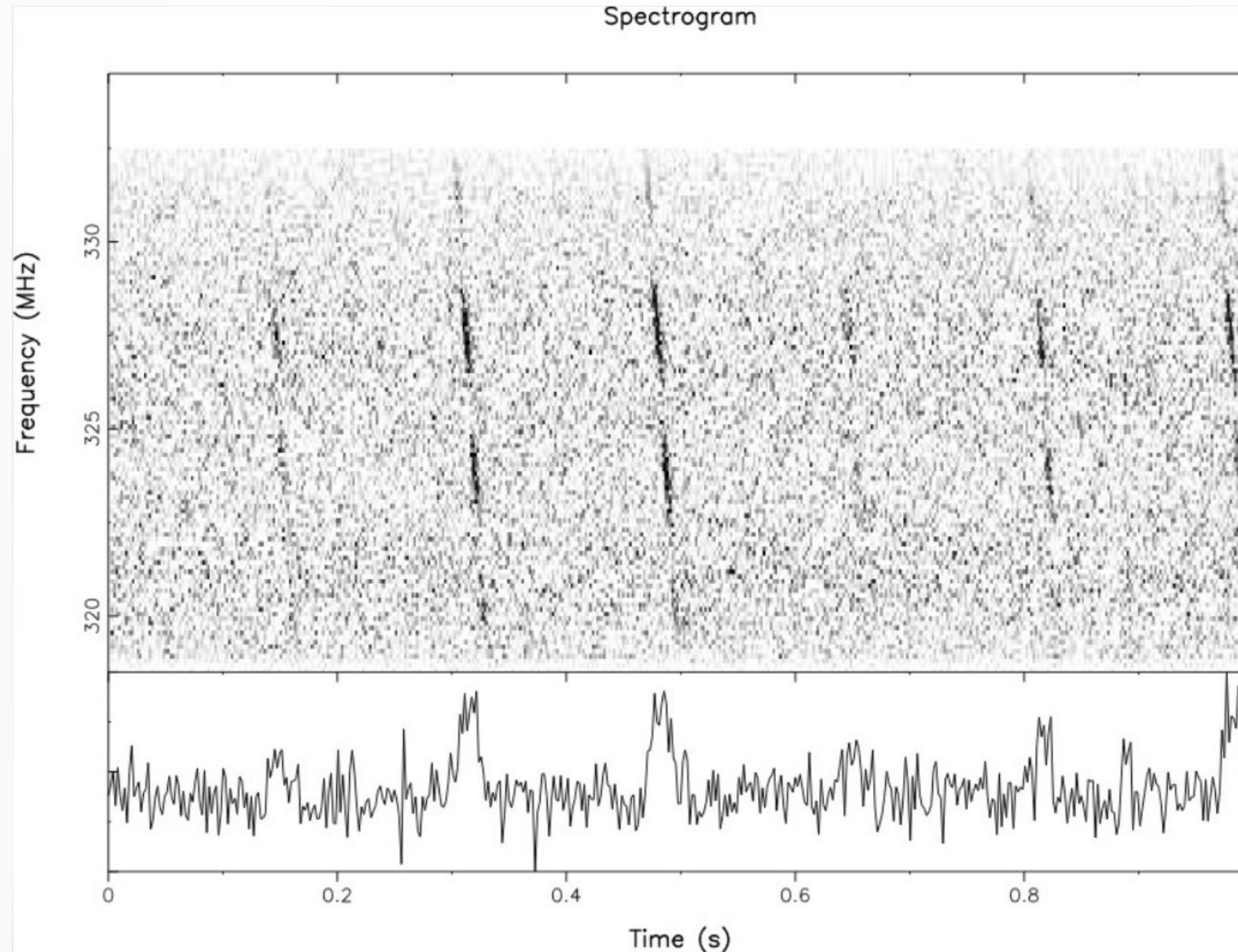
# Pulsar observations: Dedispersion



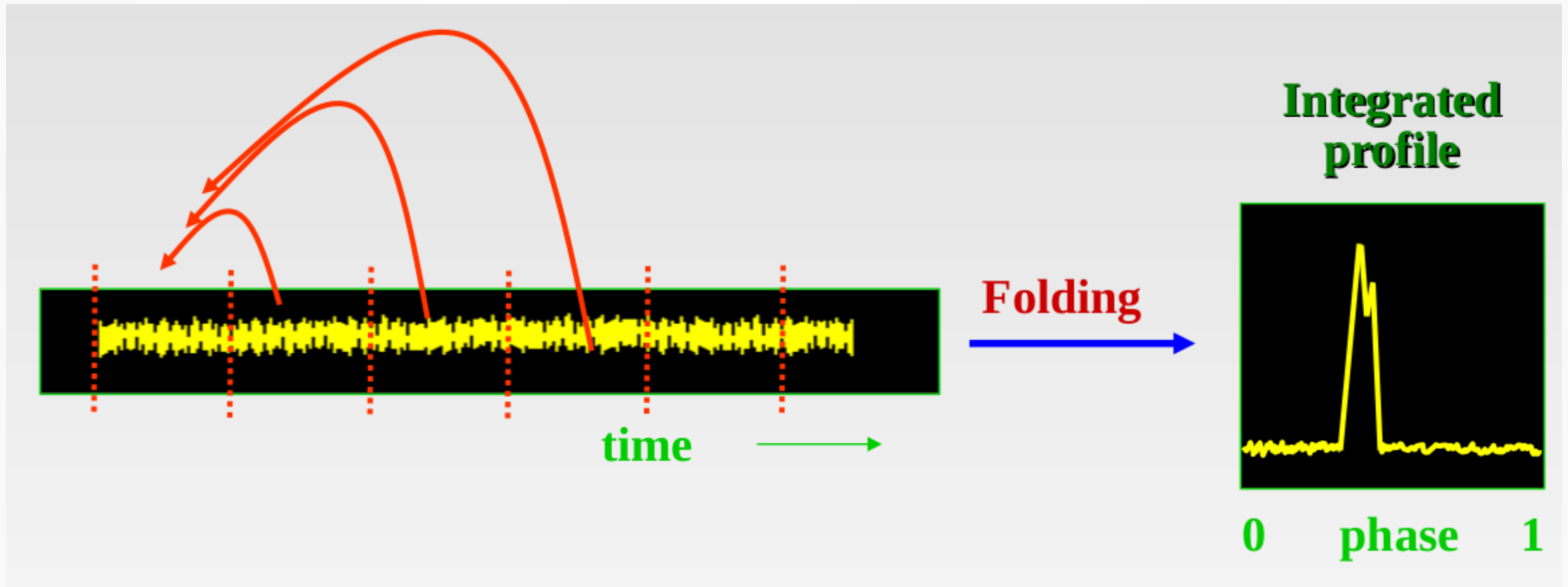
# Pulsar observations: Dedispersion



# Pulsar observations: Dedispersion



# Pulsar Observations: Folding



# Pulsar Observations: Folding

