

- Self-calibration
- Sensitivity
- W-term

Astronomical Techniques II : Lecture 13

Ruta Kale

Low Frequency Radio Astronomy (Chp. 14)

<http://www.ncra.tifr.res.in/ncra/gmrt/gmrt-users/low-frequency-radio-astronomy>

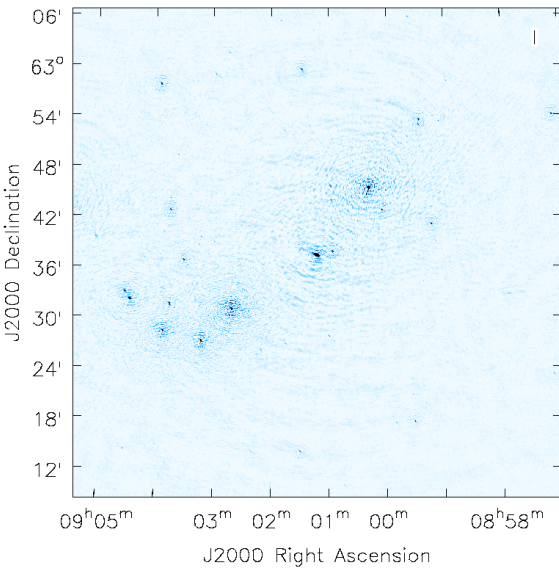
Synthesis imaging in radio astronomy II, Chp 9, 10, 19, 31, 32

Interferometry and synthesis in radio astronomy, Chp. 12

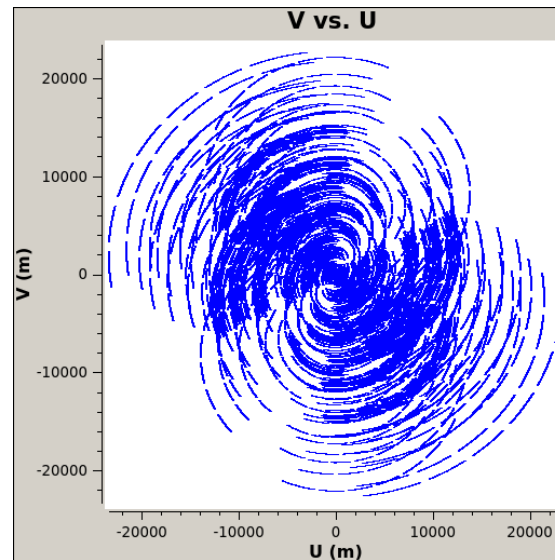
Other references are given on the slides.

Recall: Imaging

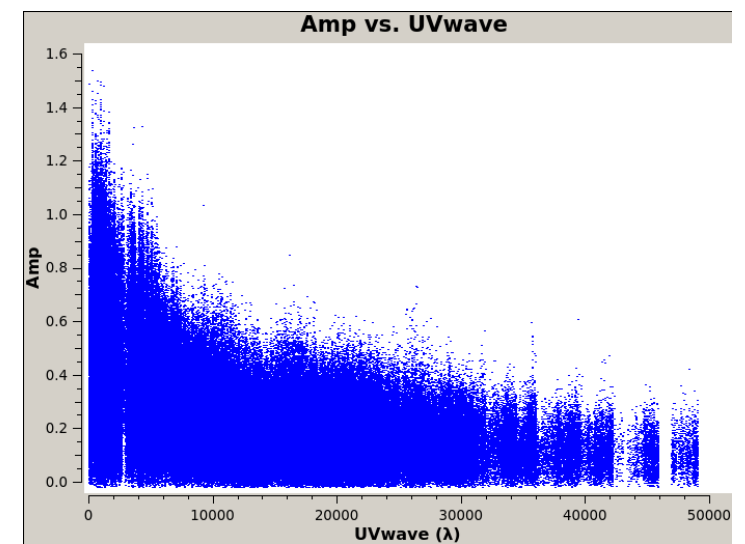
$$I^D(l, m) \equiv \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S(u, v) V'(u, v) e^{2\pi i(ul+vm)} du dv$$



Image



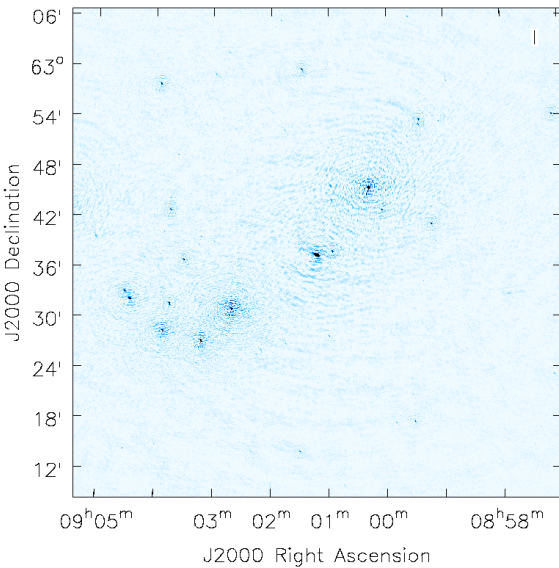
Sampling



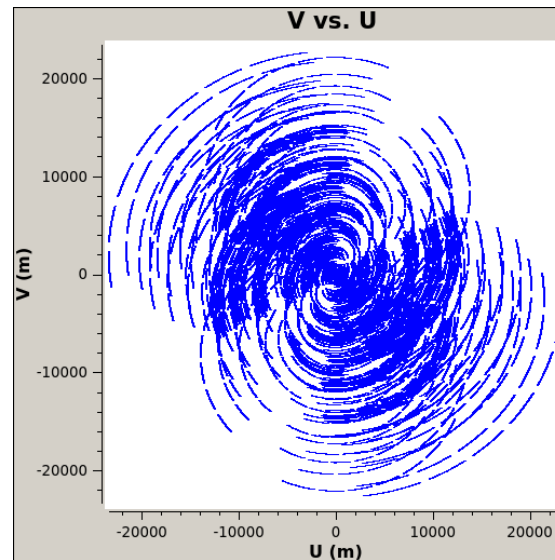
Visibilities (complex numbers)
Only amp. shown here

Recall: Imaging → Deconvolution

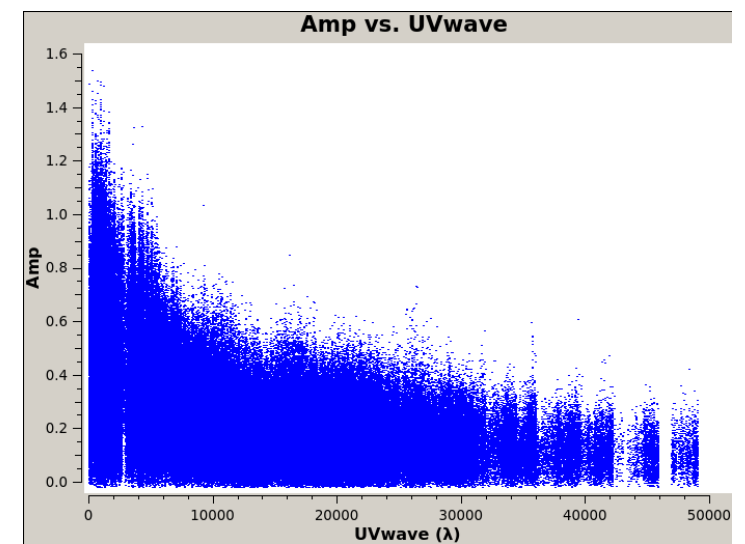
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Image



Sampling

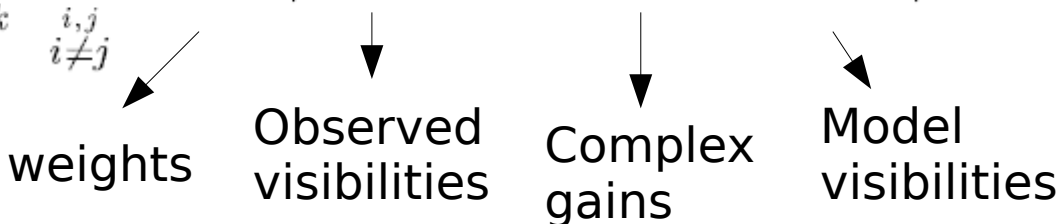


Visibilities (complex numbers)
Only amp. shown here

Self-calibration

Aim is to produce a model of the intensity distribution, the Fourier transform of which when corrected by gain factors will reproduce the measured visibilities within the noise level.

A convenient method due to Schwab 1980 is to *minimize* the sum of squares of residuals by varying complex gains g_i , g_j and the model sky,

$$\mathcal{S} = \sum_k \sum_{\substack{i,j \\ i \neq j}} w_{ij}(t_k) \left| \tilde{V}_{ij}(t_k) - g_i(t_k) g_j^*(t_k) \hat{V}_{ij}(t_k) \right|^2$$


weights Observed visibilities Complex gains Model visibilities

The time over which gains are assumed to be constant depend on the effects that govern their variation.

Self-calibration

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In most cases we have small number of degrees of freedom (the gains to be determined) and a large number of measurements of visibilities.

Self-calibration

The sky intensity model can be iteratively refined – this is done via self-calibration. The name follows after the fact that we are using the image itself as its own calibrator.

The iterative recipe is:

- Make an initial model of the sky (use CLEAN).
- Solve for complex gains.
- find the corrected visibility,

$$V_{ij,\text{corr}}(t) = \frac{\tilde{V}_{ij}(t)}{g_i(t)g_j^*(t)}$$

- Form a new model from the corrected data using constraints on the source structure.
- Again solve for complex gains and repeat until there is no improvement.

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Self-calibration is found to work. But there is no proof of convergence.

- for telescopes with large number of elements there are few variables in terms of gains as compared to the available constraints.
- sources are “simple” in their structure.

Self-calibration (caution !)

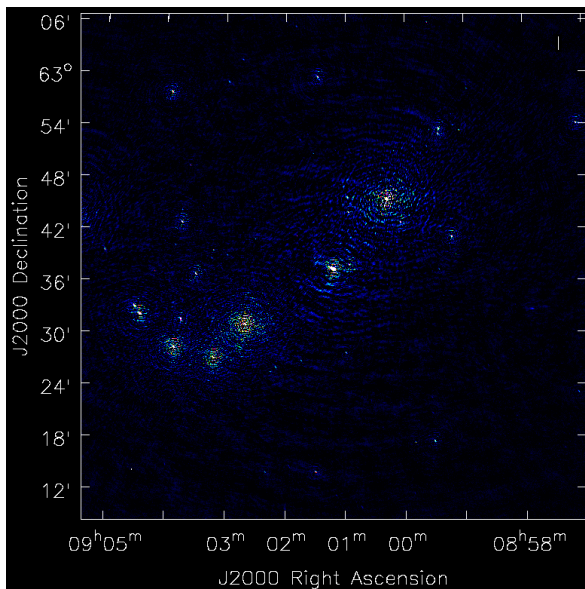
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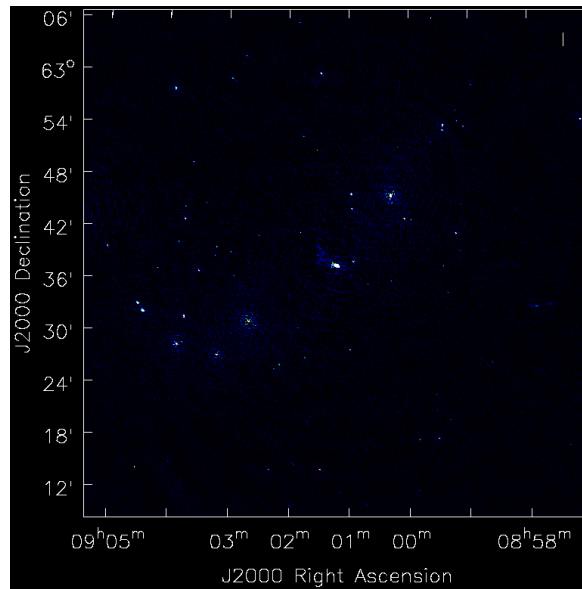
Can lead to totally wrong results if the model incorporates features from the image which are due to errors in calibration – the very effect which this procedure is trying to remove.

Solved for gain phases-only in the first few iterations and then for both amplitude and phase of the complex gain.

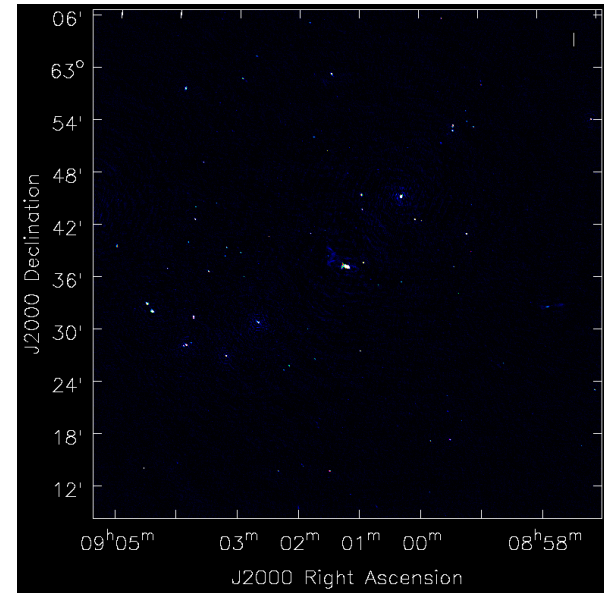
Example of self-calibration



First image after cleaning.

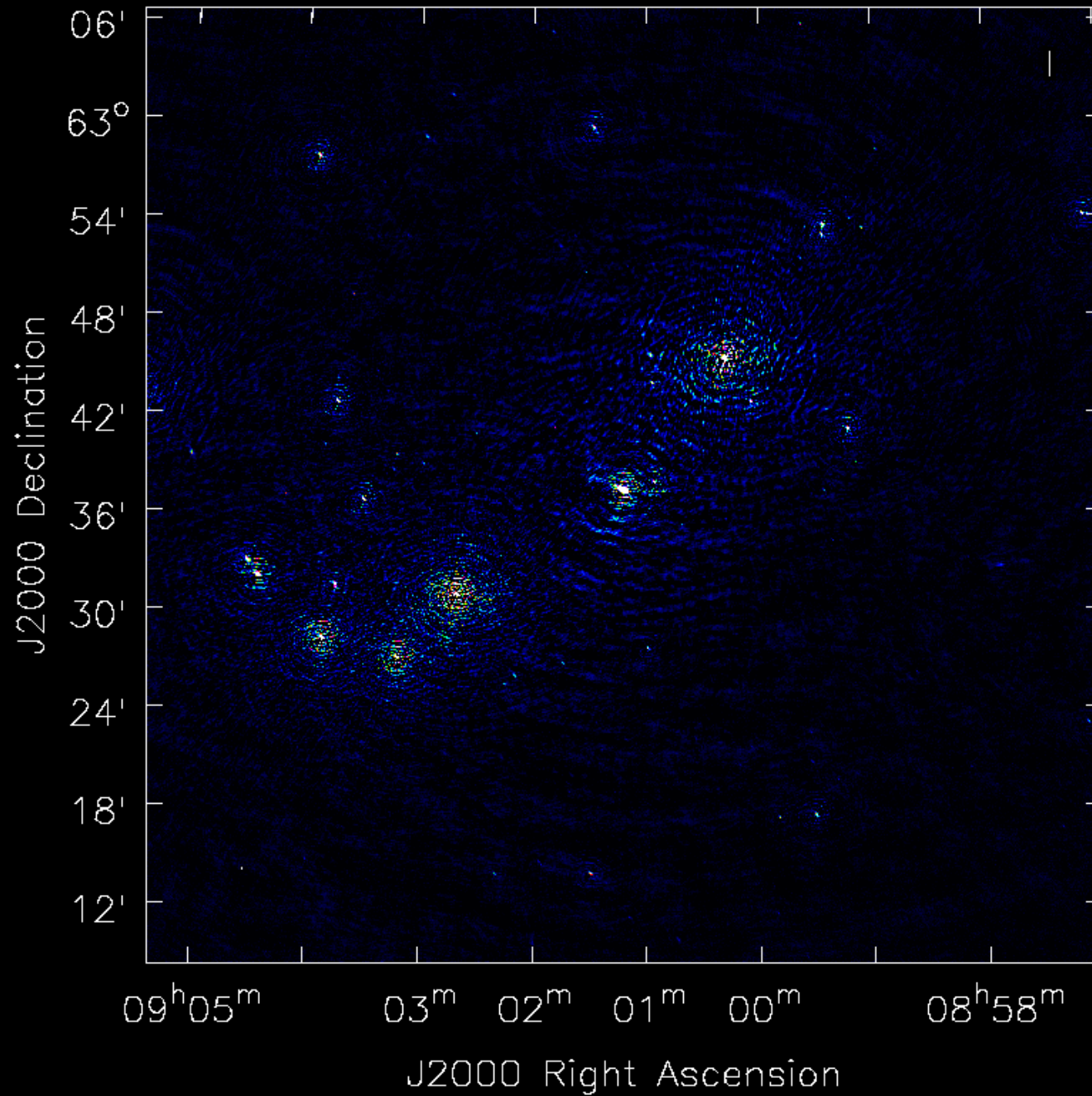


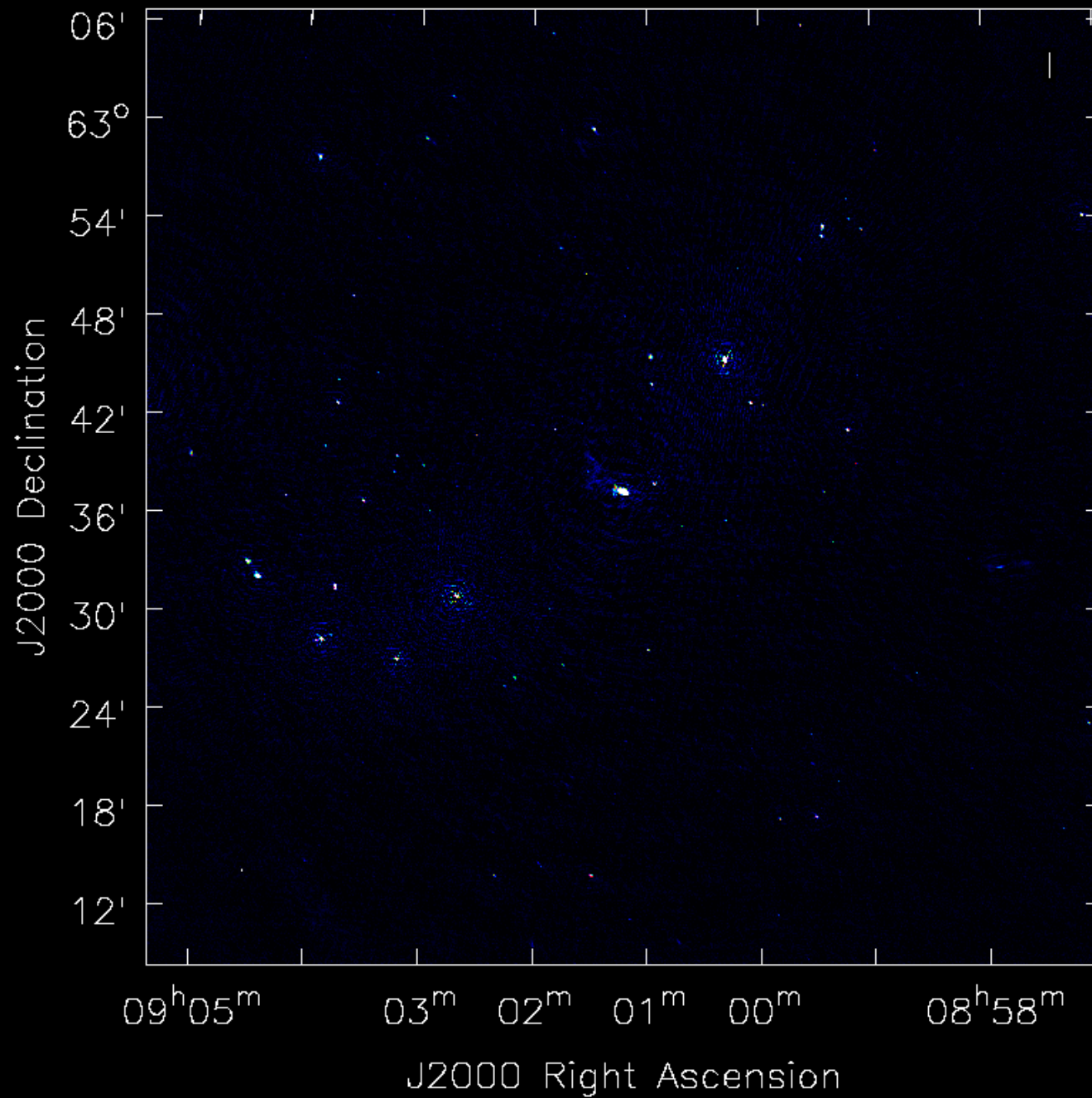
After first round of self-calibration

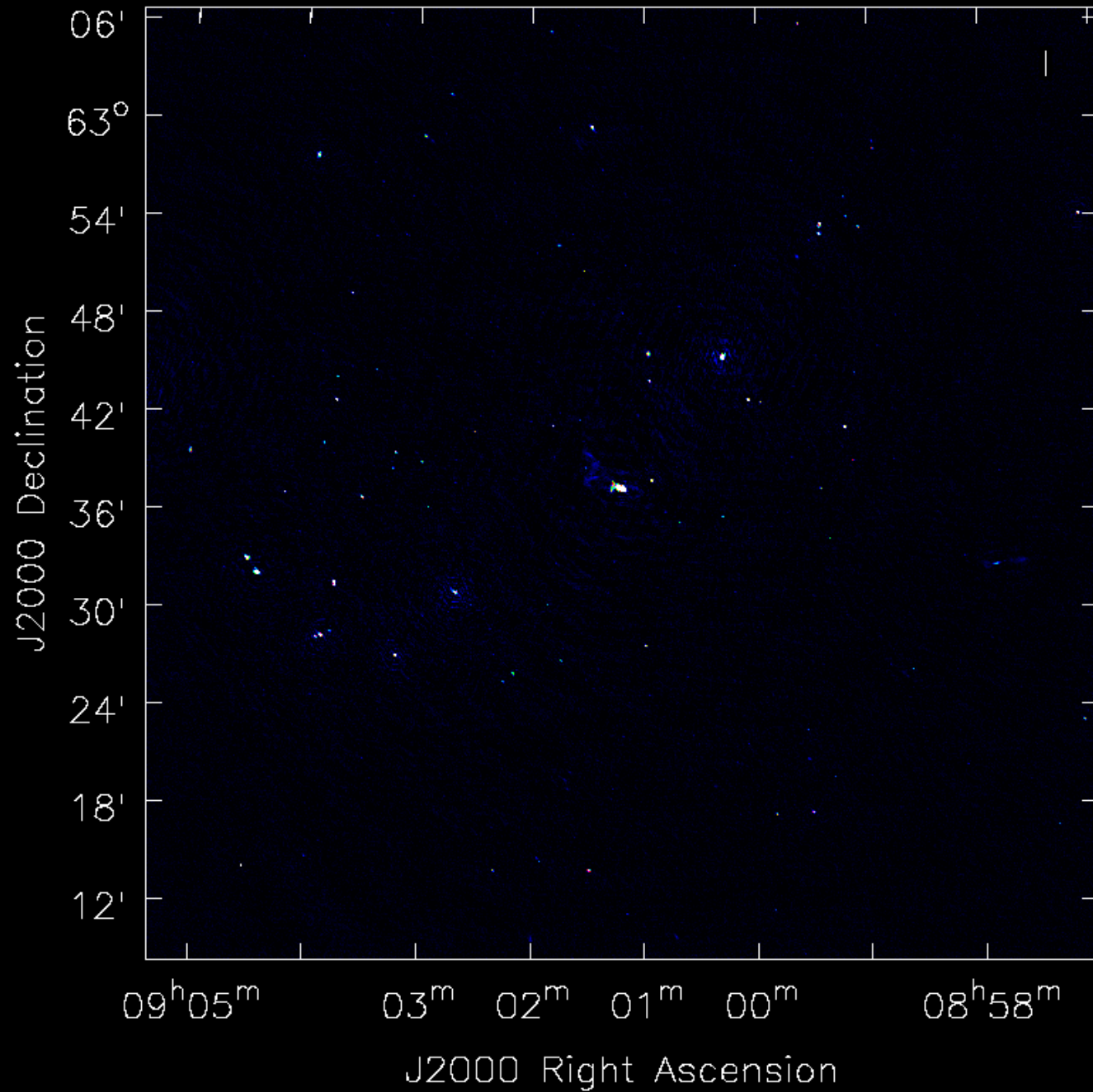


After second round of self-calibration

GMRT 610 MHz data







Closure quantities and self-calibration

$$\tilde{V}_{ij}(t) = g_i(t)g_j^*(t)V_{ij}(t) + \epsilon_{ij}(t)$$

Observed

True

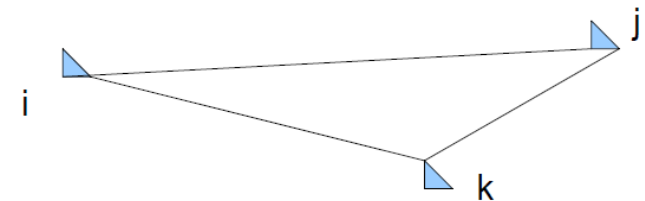


$$\tilde{\phi}_{ij}(t) = \phi_{ij}(t) + \theta_i(t) - \theta_j(t) + \text{noise term}$$

$$\theta_i(t) = \arg g_i(t)$$

Considering a loop of three antennas, the observed “closure phase” is,

$$\begin{aligned}\tilde{C}_{ijk}(t) &= \tilde{\phi}_{ij}(t) + \tilde{\phi}_{jk}(t) + \tilde{\phi}_{ki}(t) \\ &= \phi_{ij}(t) + \phi_{jk}(t) + \phi_{ki}(t) + \text{noise term} \\ &= C_{ijk}(t) + \text{noise term}\end{aligned}$$



Notice that this is independent of the individual errors, thus can be used.

Closure quantities and self-calibration

$$\tilde{V}_{ij}(t) = g_i(t)g_j^*(t)V_{ij}(t) + \epsilon_{ij}(t)$$

Observed

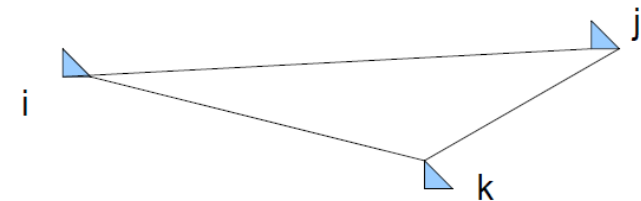
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Notice that this is independent of the individual errors, thus can be used.

For an array of N elements, there are $\frac{1}{2} N(N-1) - (N-1)$ independent closure phases - ?

A closure amplitude can be formed for a loop of 4 antennas:

$$\Gamma_{ijkl}(t) = \frac{|\tilde{V}_{ij}(t)| |\tilde{V}_{kl}(t)|}{|\tilde{V}_{ik}(t)| |\tilde{V}_{jl}(t)|}$$

Used in self-calibration by Readhead and Wilkinson 1978.

1. Make a model image.
2. for all independent closure phases, use the model to provide estimates of the true phases on two baselines and derive the phase on the other baseline in the loop from the observed phase.
3. form a new model using CLEAN from the observed visibility amplitudes and the predicted phases.
4. Repeat 2 until the model is satisfactory.

Cotton 1979 revised this scheme (least squares technique).



Sensitivity

Sensitivity is a measure of the weakest source of emission that can be detected.

In radio astronomy power is written in terms of an equivalent temperature, T , of a matched termination on the input of the receiver. Using Rayleigh-Jeans approximation to the Plancks' blackbody radiation, the power is given by:

$$P = k_B T \Delta\nu$$

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The power entering the feed is amplified by g in voltage and thus in power,

$$P_a = g^2 k_B T_a \Delta\nu \quad \text{From the source}$$

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T_a is the antenna temperature and

T_{sys} is the system temperature: includes receiver noise, feed losses, spillover, atmospheric emission, galactic background and cosmic background.

Sensitivity

$$P_a = g^2 k_B T_a \Delta\nu$$

The power from the source can be related to the flux density S , the area of the antenna A , the antenna efficiency η_a as

$$\begin{aligned} P_a &= \frac{1}{2} g^2 \eta_a A S \Delta\nu \\ &= g^2 k_B K S \Delta\nu \end{aligned}$$

Radiation received only from one channel; half for an unpolarized source.

Where $K = (\eta_a A) / (2 k_B)$

“gain” of the antenna in Kelvin of antenna temperature per Jy of flux density

K is a measure of antenna performance.

System temperature is often expressed System equivalent flux density (SEFD)
: flux density of a source that would deliver the same amount of power

$$SEFD = \frac{T_{\text{sys}}}{K} \quad \text{Jy}$$

Table 9–1. Empirical *SEFDs* from Taylor et al. (1994)

For 19 antennas
in VLBI array at
5 GHz

Taylor et al 1994

Antenna Location	Diameter (m)	<i>SEFD</i> (Jy)
NRAL, Cambridge, UK	32	140
NRAL, Jodrell Bank, UK	26	366
MPIfR, Effelsberg, Germany	100	39
OSO, Onsala, Sweden	26	757
NFRA, WSRT, Netherlands	5 × 25	133
IRA, Medicina, Italy	32	225
IRA, Noto, Italy	32	221
Haystack, Westford, MA, USA	36	606
NRAO, Green Bank, WV, USA	43	126
NRAO, VLA, NM, USA	25	319
NRAO, Saint Croix, VI, USA	25	255
NRAO, Hancock, NH, USA	25	259
NRAO, North Liberty, IA, USA	25	300
NRAO, Fort Davis, TX, USA	25	308
NRAO, Los Alamos, NM, USA	25	270
NRAO, Pie Town, NM, USA	25	280
NRAO, Kitt Peak, AZ, USA	25	308
NRAO, Owens Valley, CA, USA	25	249
NRAO, Brewster, WA, USA	25	281

Sensitivity of an interferometer

Consider an interferometer with a single real output which is the product of the voltages from the two elements. The voltage from antenna i before sampling is the sum of source voltage and noise voltage. The power from antenna i is given by, (factor a includes gain) the expectation value of the square of the voltage:

$$\begin{aligned}\langle P_i \rangle &= a_i \langle (s_i + n_i)^2 \rangle \\ &= a_i [\langle s_i^2 \rangle + \langle n_i^2 \rangle] \\ &= g_i^2 k_B (T_{ai} + T_{sysi}) \Delta\nu \\ &= g_i^2 k_B (K_i S_T + T_{sysi}) \Delta\nu\end{aligned}$$

Cross terms taken to be zero as voltages from source and noise are uncorrelated.

S_T is the total flux density seen by the antenna and we assume it is same for all the antennas.

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The power after cross multiplication in the correlator can be obtained as (i and j are the antennas):

$$\begin{aligned} \langle P_{ij} \rangle &= \frac{\sqrt{a_i a_j}}{\eta_s} \langle (s_i + n_i)(s_j + n_j) \rangle \\ &= \frac{\sqrt{a_i a_j}}{\eta_s} \langle s_i s_j \rangle \\ &= \frac{g_i g_j}{\eta_s} \sqrt{K_i K_j} k_B \Delta\nu S_c \end{aligned}$$

Efficiency factor that accounts for losses in the electronics and digital equipment

Noise from the two antennas is uncorrelated

S_c is correlated flux density

To obtain the signal-noise-ratio (SNR) we look at the RMS fluctuations of the correlator output – consider rms fluctuations of the product of the antenna voltages for each sample.

We use the fact that the square of RMS fluctuation of a Gaussian random variable (product from the correlator) is equal to the expectation value of the variable squared minus the square of the mean.

$$\sigma^2(P_{ij}) = \frac{a_i a_j}{\eta_s^2} \left\langle \left[(s_i + n_i)(s_j + n_j) \right]^2 \right\rangle - \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j k_B^2 S_c^2 \Delta\nu^2$$

Recall

$$\langle P_{ij} \rangle = \frac{\sqrt{a_i a_j}}{\eta_s} \left\langle (s_i + n_i)(s_j + n_j) \right\rangle = \frac{g_i g_j}{\eta_s} \sqrt{K_i K_j} k_B \Delta\nu S_c$$

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$$\frac{a_i a_j}{\eta_s^2} \left[2 \left\langle (s_i + n_i)(s_j + n_j) \right\rangle^2 + \left\langle (s_i + n_i)^2 \right\rangle \left\langle (s_j + n_j)^2 \right\rangle \right]$$

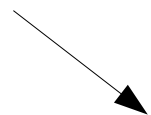
A standard relation to expand the expectation value of a product of four variables into combinations of expectation values of products of two variables is used here.

Assuming all the processes involved are Gaussian processes then we can use the properties of these processes that are already known.

If x_1, x_2, x_3 and x_4 have a joint Gaussian distribution then,

$$\langle x_1 x_2 x_3 x_4 \rangle = \langle x_1 x_2 \rangle \langle x_3 x_4 \rangle + \langle x_1 x_3 \rangle \langle x_2 x_4 \rangle + \langle x_1 x_4 \rangle \langle x_2 x_3 \rangle$$

$$\left\langle \left[(s_i + n_i)(s_j + n_j) \right]^2 \right\rangle$$



$$\left[2 \left\langle (s_i + n_i)(s_j + n_j) \right\rangle^2 + \left\langle (s_i + n_i)^2 \right\rangle \left\langle (s_j + n_j)^2 \right\rangle \right]$$

See chp 5 of LFRA.

To obtain SNR we look at the RMS fluctuations of the correlator output. We use the fact that the square of RMS fluctuation of a Gaussian random variable (product from the correlator) is equal to the expectation value of the variable squared minus the square of the mean.

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$$- \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j k_B^2 \Delta\nu^2 S_c^2$$

$$= 2 \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j (k_B \Delta\nu S_T)^2$$

$$+ \frac{g_i^2 g_j^2}{\eta_s^2} (k_B \Delta\nu)^2 (K_i S_T + T_{\text{sys}i})(K_j S_T + T_{\text{sys}j})$$

$$- \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j k_B^2 \Delta\nu^2 S_c^2$$

$$\begin{aligned} \langle P_i \rangle &= a_i \langle (s_i + n_i)^2 \rangle \\ &= g_i^2 k_B (K_i S_T + T_{\text{sys}i}) \Delta\nu \end{aligned}$$

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$$= 2 \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j (k_B \Delta\nu S_c)^2 + \frac{g_i^2 g_j^2}{\eta_s^2} (k_B \Delta\nu)^2 (K_i S_T + T_{\text{sys}i})(K_j S_T + T_{\text{sys}j}) - \frac{g_i^2 g_j^2}{\eta_s^2} K_i K_j k_B^2 \Delta\nu^2 S_c^2$$

$$\langle P_{ij} \rangle = \frac{\sqrt{a_i a_j}}{\eta_s} \left\langle (s_i + n_i)(s_j + n_j) \right\rangle = \frac{g_i g_j}{\eta_s} \sqrt{K_i K_j} k_B \Delta\nu S_c$$

$$\sigma^2(P_{ij}) = k_B^2 \Delta\nu^2 \frac{g_i^2 g_j^2}{\eta_s^2} \left(K_i K_j S_c^2 + K_i K_j S_T^2 + K_i S_T T_{\text{sys}i} + K_j S_T T_{\text{sys}j} + T_{\text{sys}i} T_{\text{sys}j} \right)$$

To obtain the noise level in units of flux density, divide by $g_i g_j \sqrt{K_i K_j} k_B \Delta\nu$

and $\sqrt{2 \Delta\nu \tau_{\text{acc}}}$

the standard deviation of the mean

↓ Source flux density to cross correlated power

$$\Delta S_{ij} = \frac{1}{\eta_s \sqrt{2 \Delta\nu \tau_{\text{acc}}}} \sqrt{S_c^2 + S_T^2 + S_T \left(\frac{T_{\text{sys}i}}{K_i} + \frac{T_{\text{sys}j}}{K_j} \right) + \frac{T_{\text{sys}i} T_{\text{sys}j}}{K_i K_j}}$$

$$\Delta S_{ij} = \frac{1}{\eta_s \sqrt{2} \Delta \nu \tau_{acc}} \sqrt{S_c^2 + S_T^2 + S_T \left(\frac{T_{sysi}}{K_i} + \frac{T_{sysj}}{K_j} \right) + \frac{T_{sysi} T_{sysj}}{K_i K_j}}$$

A square bandpass is assumed here.

For a non flat bandpass:

$$\int_0^\infty g_i(\nu) g_j(\nu) d\nu$$

$$\Delta S_{ij} = \frac{\sqrt{\int_0^\infty g_i^2(\nu) g_j^2(\nu) \left[S_c^2 + S_T^2 + S_T \left(\frac{T_{sysi}}{K_i} + \frac{T_{sysj}}{K_j} \right) + \frac{T_{sysi} T_{sysj}}{K_i K_j} \right] d\nu}}{\eta_s \sqrt{2} \tau_{acc} \int_0^\infty g_i(\nu) g_j(\nu) d\nu}$$

The noise of the correlated signal, S_c and of the source power as it adds to the total powers at each antenna, S_T , both contribute to the total noise of the correlated output.

Special cases assuming flat bandpass

$$\Delta S_{ij} = \frac{1}{\eta_s \sqrt{2} \Delta \nu \tau_{\text{acc}}} \sqrt{S_c^2 + S_T^2 + S_T \left(\frac{T_{\text{sys}i}}{K_i} + \frac{T_{\text{sys}j}}{K_j} \right) + \frac{T_{\text{sys}i} T_{\text{sys}j}}{K_i K_j}}$$

Case of strong source:

$$S_T \gg S_c$$

$$\Delta S_{ij} = \frac{S_T}{\eta_s \sqrt{2} \Delta \nu \tau_{\text{acc}}}$$

Special cases assuming flat bandpass

$$\Delta S_{ij} = \frac{1}{\eta_s \sqrt{2} \Delta\nu \tau_{\text{acc}}} \sqrt{S_c^2 + S_T^2 + S_T \left(\frac{T_{\text{sys}i}}{K_i} + \frac{T_{\text{sys}j}}{K_j} \right) + \frac{T_{\text{sys}i} T_{\text{sys}j}}{K_i K_j}}$$

Case of strong source:

$$S_T \gg S_c \quad \Delta S_{ij} = \frac{S_T}{\eta_s \sqrt{2} \Delta\nu \tau_{\text{acc}}}$$

Typically source is weak and thus ignoring terms involving flux density S:

$$\Delta S_{ij} = \frac{1}{\eta_s} \sqrt{\frac{T_{\text{sys}i} T_{\text{sys}j}}{2 \Delta\nu \tau_{\text{acc}} K_i K_j}}$$

Special cases assuming flat bandpass

$$\Delta S_{ij} = \frac{1}{\eta_s \sqrt{2} \Delta \nu \tau_{acc}} \sqrt{S_c^2 + S_T^2 + S_T \left(\frac{T_{sysi}}{K_i} + \frac{T_{sysj}}{K_j} \right) + \frac{T_{sysi} T_{sysj}}{K_i K_j}}$$

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$$S_T \gg S_c \quad \Delta S_{ij} = \frac{S_T}{\eta_s \sqrt{2} \Delta \nu \tau_{acc}}$$

Typically source is weak and thus ignoring terms involving S:

$$\Delta S_{ij} = \frac{1}{\eta_s} \sqrt{\frac{T_{sysi} T_{sysj}}{2 \Delta \nu \tau_{acc} K_i K_j}}$$

In terms of SEFD:

$$SEFD = \frac{T_{sys}}{K}$$

$$\Delta S_{ij} = \frac{1}{\eta_s} \sqrt{\frac{SEFD_i SEFD_j}{2 \Delta \nu \tau_{acc}}}$$

If SEFD same for two antennas:

$$\Delta S_{ij} = \frac{1}{\eta_s} \frac{SEFD}{\sqrt{2} \Delta \nu \tau_{acc}}$$

Sensitivity of an image

$$I_m(l, m) = C \sum_{k=1}^{2L} T_k W_k w_k V_k e^{2\pi i(u_k l + v_k m)}$$

Normalisation constant

k=0 term represents zero-spacing – not available

2L: visibility function is Hermitian – the conjugate points will also be known

Taper

Weights: Natural, uniform, etc

Weight based on S/N of individual point: important for VLBI but not for arrays with identical elements

Each image pixel is a linear combination of each measured data point.

Sensitivity of an image

$$I_m(l, m) = C \sum_{k=1}^{2L} T_k W_k w_k V_k e^{2\pi i(u_k l + v_k m)}$$

Normalisation constant

k=0 term represents zero-spacing – not available

Consider only the centre of the image:

$$I_m(0, 0) = 2 C \sum_{k=1}^L T_k W_k w_k S_{Rk} \Delta S_k$$

Taper

Weights:
Natural,
uniform,
etc

Weight based on S/N of individual point: important for VLBI but not for arrays with identical elements

Image has only real values S_R

Sensitivity of an image

$$I_m(l, m) = C \sum_{k=1}^{2L} T_k W_k w_k V_k e^{2\pi i(u_k l + v_k m)}$$

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Consider only the centre of the image:

$$I_m(0, 0) = 2 C \sum_{k=1}^L T_k W_k w_k S_{Rk}$$

Image has only real values S_R

Error in this image point is due to error ΔS_k ; thus variance in I is just the sum of variances in the Fourier components:

$$\Delta I_m = 2 C \sqrt{\sum_{k=1}^L T_k^2 W_k^2 w_k^2 \Delta S_k^2}$$

C is set to: $1 / (2 \sum_{k=1}^L T_k W_k w_k)$

To obtain the result in terms of flux density per beam area

Weight based on S/N of individual point: important for VLBI but not for arrays with identical elements

Weights: Natural, uniform, etc

Taper

2L: visibility function is Hermitian – the conjugate points will also be known

$$\Delta I_m = 2 C \sqrt{\sum_{k=1}^L T_k^2 W_k^2 w_k^2 \Delta S_k^2}$$

For no taper, natural weights, and all points of equal SNR
 $w_k = w$,

$$\begin{aligned} \Delta I_m &= 2 C \Delta S w \sqrt{\sum_{k=1}^L 1} \\ &= \frac{2 \Delta S w \sqrt{\sum_{k=1}^L 1}}{2 w \sum_{k=1}^L 1} \\ &= \Delta S \sqrt{L} / L \\ &= \Delta S / \sqrt{L} \end{aligned}$$

$$T_k = 1; W_k = 1$$

$$C = 1 / \left(2 \sum_{k=1}^L T_k W_k w_k \right)$$

$$\Delta I_m = 2 C \sqrt{\sum_{k=1}^L T_k^2 W_k^2 w_k^2 \Delta S_k^2}$$

For no taper, natural weights, and all points of weight w ,

$$\begin{aligned} \Delta I_m &= 2 C \Delta S w \sqrt{\sum_{k=1}^L 1} \\ &= \frac{2 \Delta S w \sqrt{\sum_{k=1}^L 1}}{2 w \sum_{k=1}^L 1} \\ &= \Delta S \sqrt{L} / L \\ &= \Delta S / \sqrt{L} \end{aligned}$$

N = number of antennas

$$L = \frac{1}{2} N (N - 1) (t_{\text{int}} / \tau_{\text{acc}})$$

$$\Delta I_m = \frac{1}{\eta_s} \frac{SEFD}{\sqrt{N (N - 1) \Delta \nu t_{\text{int}}}}$$

Sensitivity of a *single* polarization image formed from a homogenous array of N identical antennas:

$$\Delta I_{\text{m}} = \frac{1}{\eta_{\text{s}}} \frac{SEFD}{\sqrt{N(N-1)} \Delta\nu t_{\text{int}}}$$

For full Stokes data:

$$\Delta I = \Delta Q = \Delta U = \Delta V = \frac{\Delta I_{\text{m}}}{\sqrt{2}}$$

Sensitivity of a single polarization image formed from a homogenous array of N identical antennas:

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Noise in linear polarized flux density follows Rayleigh statistics and the position angle will follow uniform statistics.

$$P = \sqrt{Q^2 + U^2} \qquad \chi = \frac{1}{2} \tan^{-1}(U/Q)$$

See Interferometry and Synthesis in RA for more details.

Sensitivity

For a single antenna:

$$\sigma_s = \frac{2kT_{\text{sys}}}{A_{\text{eff}}(\Delta\nu_{\text{RF}}\tau)^{1/2}}$$

For an interferometer:

$$\sigma_s = \frac{2^{1/2}kT_{\text{sys}}}{A_{\text{eff}}(\Delta\nu_{\text{RF}}\tau)^{1/2}}$$

The sensitivity of an interferometer is square root 2 times better than that of each antenna but the same factor worse than a single antenna with an area of two antennas.

For an interferometer with N elements:

$$\sigma_s = \frac{2kT_{\text{sys}}}{A_{\text{eff}} [N(N-1)\Delta\nu_{\text{RF}}\tau]^{1/2}}$$

In the limit of large N, $N(N-1) \rightarrow N^2$ and the point source sensitivity of an interferometer approaches that of a single antenna whose area equals the total effective area NA_{eff} of the interferometer.

In practice, interferometers are slightly less sensitive due to the effects of sampling and digitizing.