

An aerial photograph of a lush green valley. In the foreground, a river winds through the landscape. The middle ground shows a patchwork of green fields and a small cluster of buildings. In the background, there are rolling hills under a clear blue sky with a few wispy clouds. The text is overlaid on the top half of the image.

THE INTERSTELLAR MEDIUM: III

Diffuse Atomic Gas: The HI-21cm line

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OUTLINE

- Background.
- The HI-21cm hyperfine transition.
- Absorption and emission issues.
- The spin temperature.
- HI column densities, spin and kinetic temperatures, HI masses.
- HI-21cm observational studies.

BACKGROUND

- Kinetic temperature : $f(v) = (m/2\pi kT_K)^{1/2} e^{-mv^2/2kT_K}$
- Excitation temperature : $(n_u/n_l) = (g_u/g_l) e^{-hv/kT_X}$
- Radiation temperature : $B(v, T_R) = (2hv^3/c^2) [e^{hv/kT_R} - 1]^{-1}$

Rayleigh-Jeans limit : $B(\lambda, T) = (2kT/\lambda^2)$

- Einstein A-coefficient : $A_{ul} = (64\pi^4 e^2 v^3 / 3hc^3) (g_u/g_l) |D_{ul}|^2$

- Equation of radiative transfer:

$$I_\nu(\tau_\nu) = I_\nu(0) e^{-\tau_\nu} + \int_0^{\tau_\nu} e^{-(\tau_\nu - \tau)} S_\nu d\tau \quad , \quad \text{where } S_\nu \equiv (j_\nu/\kappa_\nu).$$

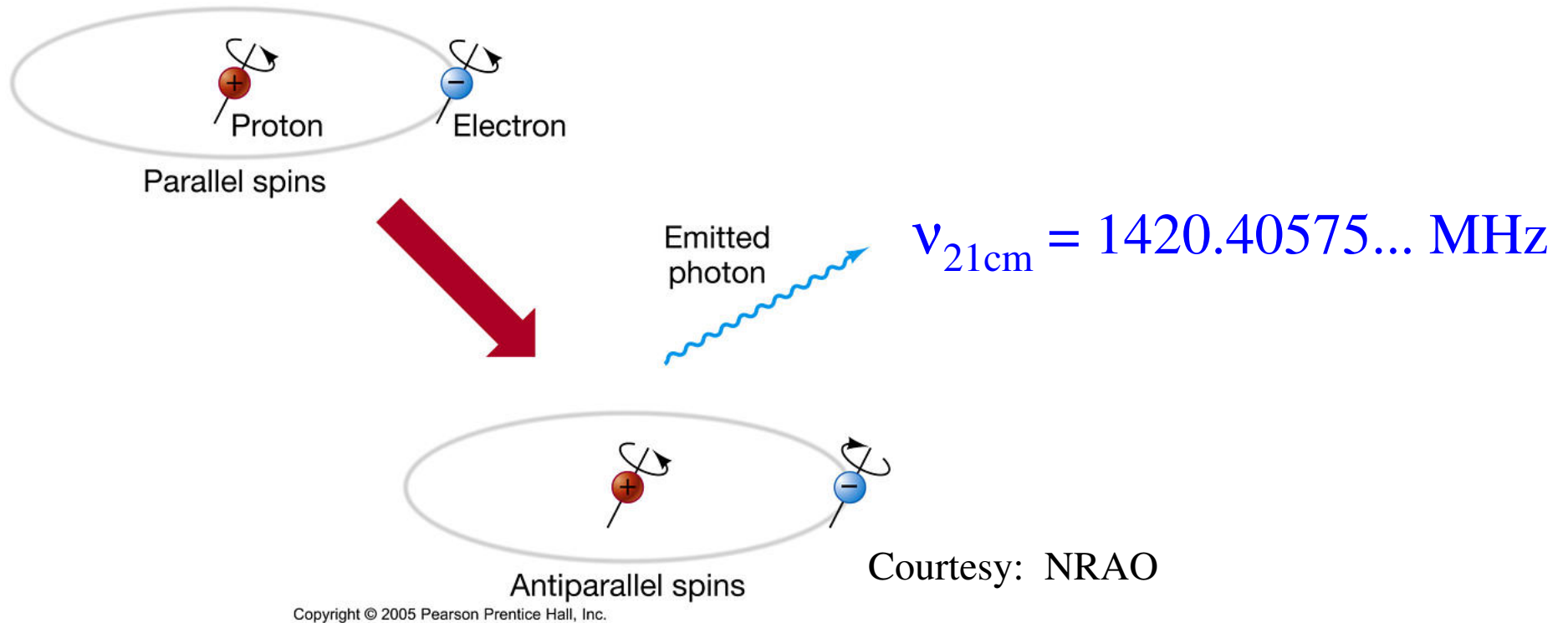
$j_\nu \equiv$ Emissivity; Attenuation coefficient, $\kappa_\nu = n_l \sigma_{lu}(\nu) - n_u \sigma_{ul}(\nu)$.

- For a uniform medium, with level populations given by T_X :

$$I_\nu = I_\nu(0) e^{-\tau_\nu} + B_\nu(T_X) (1 - e^{-\tau_\nu})$$

THE HI-21CM LINE

- “Spin-flip” transition: e^-p^+ spins go parallel to anti-parallel.



- “Forbidden” magnetic dipole transition: $A_{21\text{cm}} = 2.87 \times 10^{-15} \text{ s}^{-1}$.
- Large mean lifetime $\sim 1/A_{21\text{cm}} \approx 10^7$ years!
- Narrow natural broadening, $\Delta\nu_{21\text{cm}} \sim A_{21\text{cm}}/2\pi \approx 5 \times 10^{-16} \text{ Hz}$!

THE HI-21CM LINE: ABSORPTION ISSUES

- Level populations: $(n_1/n_0) = (g_1/g_0) e^{-hv/kT_S}$.
 $T_S \equiv$ “Spin temperature”; $g = (2S + 1)$; $g_1 = 3$, $g_0 = 1$.
- $hv/k \approx 0.07$ K, $\ll T_S \Rightarrow (n_1/n_0) \approx 3 (1 - hv/kT_S) \Rightarrow n \approx 4n_0$.
- The absorption cross-section is given by
$$\sigma_{10}(v) = (g_1/g_0) (c^2/8\pi v^2) A_{21\text{cm}} \phi(v).$$
- The attenuation coefficient, $\kappa_v = n_0 \sigma_{01}(v) - n_1 \sigma_{10}(v)$.
$$\begin{aligned} \Rightarrow \kappa_v &= (c^2/8\pi v^2) \cdot 3 \cdot A_{21\text{cm}} \cdot (n/4) \cdot \phi(v) [1 - e^{-hv/kT_S}] \\ &= (3c^2/32\pi v^2) \cdot A_{21\text{cm}} \cdot n \cdot \phi(v) (hv/kT_S) \end{aligned}$$
- Optical depth, $\tau_v = \int \kappa_v ds = (3c^2/32\pi v^2) A_{21\text{cm}} \phi(v) \cdot (hv/kT_S) N_{\text{HI}}$
 \Rightarrow Total HI column density, $N_{\text{HI}} = 1.82 \times 10^{18} \int T_S \tau_v dV$.

THE HI-21CM LINE: EMISSION ISSUES

- Level populations: $(n_1/n_0) = (g_1/g_0) e^{-hv/kT_S}$.
- $hv/k \approx 0.07 \text{ K} (\ll T_S) \Rightarrow (n_1/n_0) \approx 3 (1 - hv/kT_S) \Rightarrow n \approx 4n_0$.

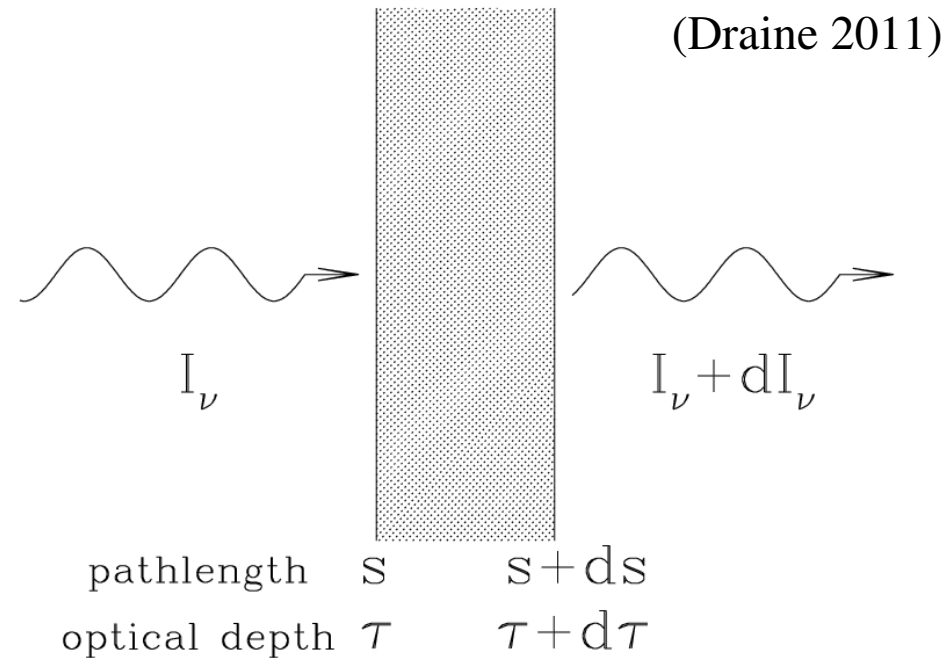
- Equation of radiative transfer:

$$dI_\nu = -I_\nu \kappa_\nu ds + j_\nu ds$$

- $I_\nu = I_\nu(0) e^{-\tau_\nu} + \int_0^{\tau_\nu} e^{-(\tau_\nu - \tau)} (j_\nu/\kappa_\nu) d\tau$.

- Kirchhoff's law: $(j_\nu/\kappa_\nu) \equiv B_\nu(T_S)$

$$\Rightarrow I_\nu = I_\nu(0) e^{-\tau_\nu} + B_\nu(T_S) (1 - e^{-\tau_\nu}).$$



- For HI-21cm, $(hv/kT_S) \ll 1 \Rightarrow T_B = T_B(0) e^{-\tau_\nu} + T_S (1 - e^{-\tau_\nu})$.

- If there's no background source $\Rightarrow T_B = T_S (1 - e^{-\tau_\nu})$.

HI-21CM STUDIES

- The HI column density, $N_{\text{HI}} = 1.8 \times 10^{18} \int T_S \tau_v dV$.
- Isothermal cloud: The brightness temperature, $T_B = T_S (1 - e^{-\tau_v})$.
- For optically-thin absorption, $\tau_v \ll 1 \Rightarrow T_B = T_S \cdot \tau_v$
 $\Rightarrow N_{\text{HI}} = 1.8 \times 10^{18} \int T_B dV$
 \Rightarrow Can measure N_{HI} directly from HI-21cm emission studies!
- In the general case, $N_{\text{HI}} = 1.8 \times 10^{18} \int T_B \tau_v [1 - e^{-\tau_v}]^{-1} dV$.
- HI-21cm optical depth depends on *both* N_{HI} & T_S ; $\tau_v \propto (1/T_S)$
 \Rightarrow More difficult to detect absorption from warm HI.
- Absorption by an isothermal cloud: $N_{\text{HI}} = 1.8 \times 10^{18} T_S \int \tau_v dV$.
For multiple clouds on the sightline: $N_{\text{HI}} = 1.8 \times 10^{18} \langle T_S \rangle \int \tau_v dV$;
Column-density-weighted harmonic mean $\langle T_S \rangle = [\sum_i (n_i/T_i)]^{-1}$

THE SPIN TEMPERATURE

- Level populations determined by the kinetic temperature, the HI-21cm radiation field and the Lyman- α colour temperature.

- For collisional excitation or de-excitation (i.e. ignoring radiation):

$$(dn_1/dt) = n_c n_0 k_{01} - n_c n_1 k_{10} - n_1 A_{10}$$

k_{01} & k_{10} are collisional rate coefficients: $k_{01} = (g_1/g_0) k_{10} e^{-hv/kT_K}$

$n_c \equiv$ density of the colliding partner.

In the steady state, $(dn_1/dt) = 0 \Rightarrow (n_1/n_0) = n_c k_{01} / [n_c k_{10} + A_{10}]$

- Including a radiation field, of specific energy density u_ν

$$\Rightarrow (dn_1/dt) = n_0 [n_c k_{01} + n_\gamma (g_1/g_0) A_{10}] - n_1 [n_c k_{10} + (1 + n_\gamma) A_{10}]$$

where n_γ is the photon occupancy number = $(c^3/8\pi h\nu^3)u_\nu$

- For black-body radiation, $n_\gamma = [e^{hv/kT_R} - 1]^{-1}$.

- Steady state: $(n_1/n_0) = [n_c k_{01} + n_\gamma (g_1/g_0) A_{10}] / [n_c k_{10} + (1 + n_\gamma) A_{10}]$

THE SPIN TEMPERATURE

- $(n_1/n_0) = [n_c k_{01} + n_\gamma (g_1/g_0) A_{10}] / [n_c k_{10} + (1 + n_\gamma) A_{10}]$
- **Critical density:** Collisional de-excitation = Radiative de-excitation
$$n_{crit,u} = [\sum_{l<u} (1 + n_\gamma) A_{ul}] / [\sum_{l<u} k_{ul}]$$

If $n_c \gg n_{crit,u} \Rightarrow T_X = T_K$; if $n_c \ll n_{crit,u} \Rightarrow T_X = T_R$.
- **Note:** Black-body radiation & $T_R = T_K \Rightarrow T_X = T_K$ for *all* densities!
- For the HI-21cm line, $n_{crit} = (1 + n_\gamma) A_{10} / k_{10}$.

Main collision partners: HI atoms. (e.g. Allison & Dalgarno 1969)

Radiation field temperature $\sim (2.73 + 1.04) \text{ K} = 3.77 \text{ K}$

\Rightarrow Photon occupancy number $\sim (3.77 / 0.07) \sim 55$

$\Rightarrow n_{crit} = 1.7 \times 10^{-3} (T_K/100 \text{ K})^{-0.66} \text{ cm}^{-3}$ (50 K \leq T \leq 200 K)
- Solve $T_S = 0.07 / \ln(n_0 g_1 / n_1 g_0)$ for different n_c and T_K values.

THE SPIN TEMPERATURE

- $(n_1/n_0) = [n_c k_{01} + n_\gamma (g_1/g_0) A_{10}] / [n_c k_{10} + (1 + n_\gamma) A_{10}]$

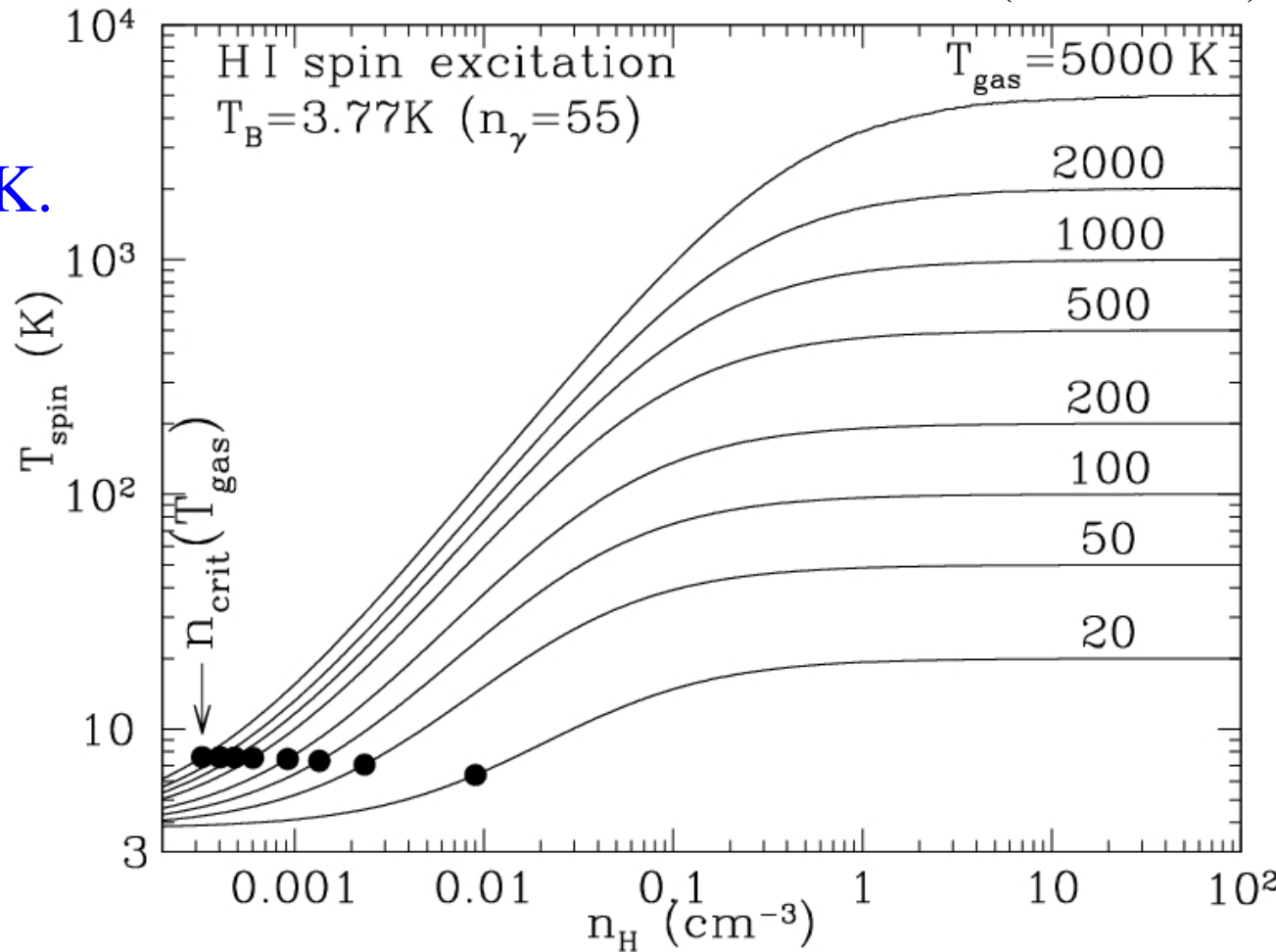
(Draine 2011)

- If $n_c \gg n_{crit} \Rightarrow T_S = T_K$.

- If $n_c \ll n_{crit} \Rightarrow T_S = 3.77 \text{ K}$.

- $n_c > 1 \text{ cm}^{-3} \Rightarrow T_S \approx T_K$.

- $n_c < 1 \text{ cm}^{-3} \Rightarrow T_S < T_K$!



- Collisions insufficient to thermalize HI-21cm levels at low densities.

- Wouthysen-Field mechanism: Lyman- α radiation field ?

- Also insufficient for typical Lyman- α intensities.

(Liszt 2001)

HI-21CM STUDIES: OBSERVABLES

- Emission studies: If $\tau_v \ll 1 \Rightarrow N_{\text{HI}} = 1.8 \times 10^{18} \int T_B dV$
 \Rightarrow Can measure N_{HI} directly from HI-21cm emission studies!
- Absorption studies : $N_{\text{HI}} = 1.8 \times 10^{18} T_S \int \tau_v dV$.
 \Rightarrow Can infer T_S if N_{HI} known (HI-21cm emission or Ly- α).
- If $\tau_v \ll 1$, Gaussian line profile in local thermal equilibrium.
 \Rightarrow Can fit multi-Gaussian profile to infer kinetic temperature.
- Zeeman splitting: Shift of 2.8 Hz / μG between 2 polarizations.
 \Rightarrow Infer parallel magnetic field via polarization studies.
- External galaxies: HI mass, $M_{\text{HI}} = 2.35 \times 10^5 D^2 \int S dV$ (M_\odot)
Mapping studies: Gas mass, distribution, velocity field.