# Astronomical Techniques II <br> Lecture 13 - Some things not covered in this course 

Divya Oberoi<br>IUCAA NCRA Graduate School div@ncra.tifr.res.in<br>March-May 2014

## A formalism for 3-D imaging

1

$$
V(u, v, w)=\int_{e_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\mathcal{A}(I, m) I(I, m)}{\sqrt{1-I^{2}-m^{2}}}}^{e^{-2 \pi i\left(u l+v m+w\left(\sqrt{1-l^{2}-m^{2}}-1\right)\right)} d l d m}
$$

2

$$
\begin{aligned}
V(u, v, w) e^{-2 \pi i w}= & \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\mathcal{A}(I, m) I(I, m)}{\sqrt{1-I^{2}-m^{2}}} \\
& \delta\left(n-\sqrt{1-I^{2}-m^{2}}\right) \\
& e^{-2 \pi i(u l+v m+w n)} d l^{\infty} d m d n
\end{aligned}
$$

1

$$
\begin{aligned}
& I^{D(3)}= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} v(u, v, w) S(u, v, w) e^{-2 \pi i w} \\
& e^{2 \pi i(u l+v m+w n)} d u d v d w
\end{aligned}
$$

$2 I^{D(3)}=I^{(3)} \star B^{D(3)}$ where,

$$
I^{(3)}(I, m, n)=\frac{\mathcal{A}(I, m) I(I, m)}{\sqrt{1-I^{2}-m^{2}}} \delta\left(n-\sqrt{1-I^{2}-m^{2}}\right)
$$

## 3D Imaging



Figure 19-1. The image volume and its relation to the sky brightness. (Left) Threedimensional transformation of the analytic visibility function maps the sky brightness onto a unit sphere. The dots represent these sources. (Middle) Convolution with a dirty beam results in sidelobes, shown as dashed lines, throughout the volume above and below the unit sphere. (Right) After deconvolution, the images are represented by finite-size "clean beams" on the unit sphere. The two-dimensional image is recovered by projection onto the tangent plane, indicated by vertical dashed lines.

## 3D Imaging



Figure 19-2. The image volume and its relation to a 'standard' two-dimensional image. (Left) At a particular time, a 'snapshot' with a two-dimensional array will project the true structure on the unit sphere onto the tangent plane with a 'ray beam', tilted at a particular angle given by the geometry of the array at the time of observation. (Right) At a later time, the array geometry has changed due to earth rotation, so the projection is now at a different angle. The apparent positions of the objects which are not located at the tangent point have changed with respect to the earlier observation.

## 3D Imaging

1 Faceting/Polyhedron imaging
1 Divide the image into many many facets, each small enough that the small FoV and small $w$ term approximation are sastified within it
2 CLEAN flux is subtracted from ungridded visibilities
3 No. of facets depends upon FoV and resolution
4 100-1000 times slower than 2D imaging
2 w projection (Cornwell, Golap and Bhatnagar, 2008)
$1 V(u, v, w)=$
$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{I(I, m)}{\sqrt{1-l^{2}-m^{2}}} \mathbf{G}(I, m, w) e^{-2 \pi i(u l+v m)} d l d m$,
where

$$
\mathbf{G}(I, m, w)=e^{-2 \pi i\left(w \sqrt{1-I^{2}-m^{2}}-1\right)}
$$

$2 V(u, v, w)=\hat{G}(u, v, w) \star V(u, v, w=0)$
3 Order of magnitude faster

## Polarization

1 Most non-thermal processes give rise to at least partially polarised emission
2 Polarized emission is an important diagnostic of the conditions in the radio source and in the intervening medium

3
4 DoF needed to describe the polarization state of radiation
5 A given feed is sensitive to only one of the orthogonal pols (linear or circular)
6 Measure both polarizations and compute all four cross-correlations

## Polarization Measurements

1 Significantly harder than total intensity
1 For the vast majority of sources, fractional polarization is quite low - pushed into low SNR regimes
2 Number of DoF for imaging increase by a factor of 4 - less well constrained
$[3$ Calibration issues
1 Instrumental
2 Propagation
3 Need for polarization calibrator
4 calibration has a strong direction dependence (absolute, as well as within the fov) 1
5 Alt-Az mounts

## The Hamaker-Bregman-Sault Measurement Equation

I Hamaker, Bregman and Sault - 1996-1998
2 Jones Matrix
$1 E_{0} \cos (\omega t+\phi)=E_{0} e^{i \phi}$
$\boldsymbol{2}\binom{E_{R}^{\prime}}{E_{L}^{\prime}}=\left(\begin{array}{ll}a & b \\ c & d\end{array}\right)\binom{E_{R}}{E_{L}}$
(3) $J_{\text {gain }}=\left(\begin{array}{cc}g_{R} & 0 \\ 0 & g_{L}\end{array}\right)$

4 J $J_{\text {leakage }}=\left(\begin{array}{cc}1 & D_{R} \\ -D_{L} & 1\end{array}\right)$
5 J $J_{\text {rotation }}=\left(\begin{array}{cc}e^{-i \theta} & 0 \\ 0 & e^{i \theta}\end{array}\right)$
6 $J_{\text {overall }}=J_{\text {gain }} J_{\text {leakage }} J_{\text {rotation }} \ldots$

## Jones matrices

11 Js are different for each antenna, and are usually time and frequency dependent

2 Provide a framework to represent propagation of signal path up to the correlator
3 Complicated systems can be handled gracefully
4 Provides an approach which allows individual effects to be modelled in different physically relevant manners
5 Matrix formulation - well suited for computational scalability and efficiency

## Jone matrices - Polarimetric Equivalent

(1) $V_{i, j}^{\prime}=g_{i} g_{j}^{*} V_{i, j}$
(2 $\mathbf{A} \otimes \mathbf{B}=a_{i, j} \mathbf{B}$

$$
\left(\mathbf{A}_{i} \mathbf{B}_{i}\right) \otimes\left(\mathbf{A}_{j} \mathbf{B}_{j}\right)=\left(\mathbf{A}_{i} \otimes \mathbf{A}_{j}\right)\left(\mathbf{B}_{i} \otimes \mathbf{B}_{j}\right)
$$

3 Inputs to the correlator $-E_{i}^{\prime}=\mathbf{J}_{i} E_{i}$
4 Outputs of the correlator $-E_{i}^{\prime} \otimes E_{j}^{\prime *}$

$$
\left(\mathbf{J}_{i} E_{i}\right) \otimes\left(\mathbf{J}_{j} E_{j}\right)^{*}=\left(\mathbf{J}_{i} \otimes \mathbf{J}_{j}^{*}\right)\left(E_{i} \otimes E_{j}^{*}\right)
$$

$\boldsymbol{5} E_{i}^{\prime} \otimes E_{j}^{\prime *}=\left(\begin{array}{c}E_{R, i} E_{R, j}^{*} \\ E_{R, i} E_{L, j}^{*} \\ E_{L, j} E_{R, j}^{*} \\ E_{L, i} E_{L, j}^{*}\end{array}\right)$
$\boldsymbol{1}\left\langle E_{i}^{\prime} \otimes E_{j}^{\prime *}\right\rangle=\left(\begin{array}{l}V_{R R, i j} \\ V_{R L, j j} \\ V_{L R, i j} \\ V_{L L, i j}\end{array}\right)$
[ $V_{i j}^{\prime}=\left(\mathbf{J}_{i} \otimes \mathbf{J}_{j}^{*}\right) V_{i j}-V_{i j}$ - Coherency vector
3 Calibration requires estimating the different $J_{i}$ s and applying the inverse matrix to the measured Coherency vector

## Relationship with Stokes vectors

$\boldsymbol{1} V_{S}=\left(\begin{array}{c}V_{l} \\ V_{Q} \\ V_{U} \\ V_{V}\end{array}\right)$ - Stokes Visibility Vector
$2 V_{i j}=\mathbf{S} V_{S, i j}$
$3 V_{i j}^{\prime}=\left(\mathbf{J}_{i} \otimes \mathbf{J}_{j}\right) \mathbf{S} V_{S, i j}$
$4 \mathbf{S}_{\text {circ }}=\left(\begin{array}{cccc}1 & 0 & 0 & 1 \\ 0 & 1 & i & 0 \\ 0 & 1 & -i & 0 \\ 1 & 0 & 0 & -1\end{array}\right) \quad \mathbf{S}_{\text {linear }}=\left(\begin{array}{cccc}1 & 1 & 0 & 0 \\ 0 & 0 & 1 & i \\ 0 & 0 & 1 & -i \\ 1 & -1 & 0 & 0\end{array}\right)$

## Direction Dependent Effects

$11 V_{i j}^{\prime}=$
$\left(\mathbf{J}_{V I S, i} \otimes \mathbf{J}_{V I S, j}^{*}\right) \sum_{k}\left(\mathbf{J}_{S K Y, i}\left(\theta_{k}, \phi_{k}\right) \otimes \mathbf{J}_{S K Y, j}^{*}\left(\theta_{k}, \phi_{k}\right)\right) \mathbf{S} V_{S, i j, k}$
$[2$ Assumptions
1 Sky can be represented as a colletion of point sources
2 Position dependent effects preceed position independent effects

3 Jones matrices formalism - not valid in presence of depolarization (analog of loss of coherence)
4 Generalized to include this via Mueller matrices

## Propagation Effects

1 Media
11 Interplanetary Medium
$[2$ lonosphere
2 Impacts
1 Refraction
2 Absorption
3 Change in polarization properties of the radiation

## Ionospheric Propagation Effects



## Ionospheric Propagation Effects



## Spectral Line Observations

1 SNR considerations for a single spectral channel rather than the full bandwidth
2 Spectral baseline subtraction

## VLBI

## 1 Highest resolution observations at any wavelength $60 \mu$ arcsec <br> 2 About the size of an orange on the Moon

Table 1-1. Terms of a VLBI Geometric Model ${ }^{a}$

| Item | Approx max Magnitude ${ }^{b}$ | Time scale |
| :--- | ---: | ---: |
| Zero order geometry. | 6000 km | 1 day |
| Nutation | $\sim 20^{\prime \prime}$ | $<18.6 \mathrm{yr}$ |
| Precession | $\sim 0.5$ arcmin $/ \mathrm{yr}$ | years |
| Annual aberration | $20^{\prime \prime}$ | 1 year |
| Retarded baseline | 20 m | 1 day |
| Gravitational delay | 4 mas $@ 90^{\circ}$ from sun | 1 year |
| Tectonic motion | $10 \mathrm{~cm} / \mathrm{yr}$ | years |
| Solid Earth Tide | 50 cm | 12 hr |
| Pole Tide | 2 cm | $\sim 1 \mathrm{yr}$ |
| Ocean Loading | 2 cm | 12 hr |
| Atmospheric Loading | 2 cm | weeks |
| Post-glacial Rebound | several mm/yr | years |
| Polar motion | $0.5^{\prime \prime}$ | $\sim 1.2$ years |
| UT1 (Earth rotation) | Random at several mas | Various |
| Ionosphere | $\sim 2 \mathrm{mat} 2 \mathrm{GHz}$ | seconds to years |
| Dry Troposphere | 2.3 m at zenith | hours to days |
| Wet Troposphere | $0-30 \mathrm{~cm}$ at zenith | seconds to seasonal |
| Antenna structure | $<10 \mathrm{~m} .1 \mathrm{~cm}$ thermal | - |
| Parallactic angle | 0.5 turn | hours |
| Station clocks | few microsec | hours |
| Source structure | 5 cm | years |

## VLBI - Looking for the Black Hole at the center of our Galaxy


1.3 mm VLBI - SMT, JCMT, CARMA The solid curve is the circular Gaussian fit with FWHM of 43 micro arcseconds. Dashed line is for annular ring model of inner radius 35 micro arcseconds and outer radius of 80 micro arcseconds (Doeleman et al 2008, Nature)

