

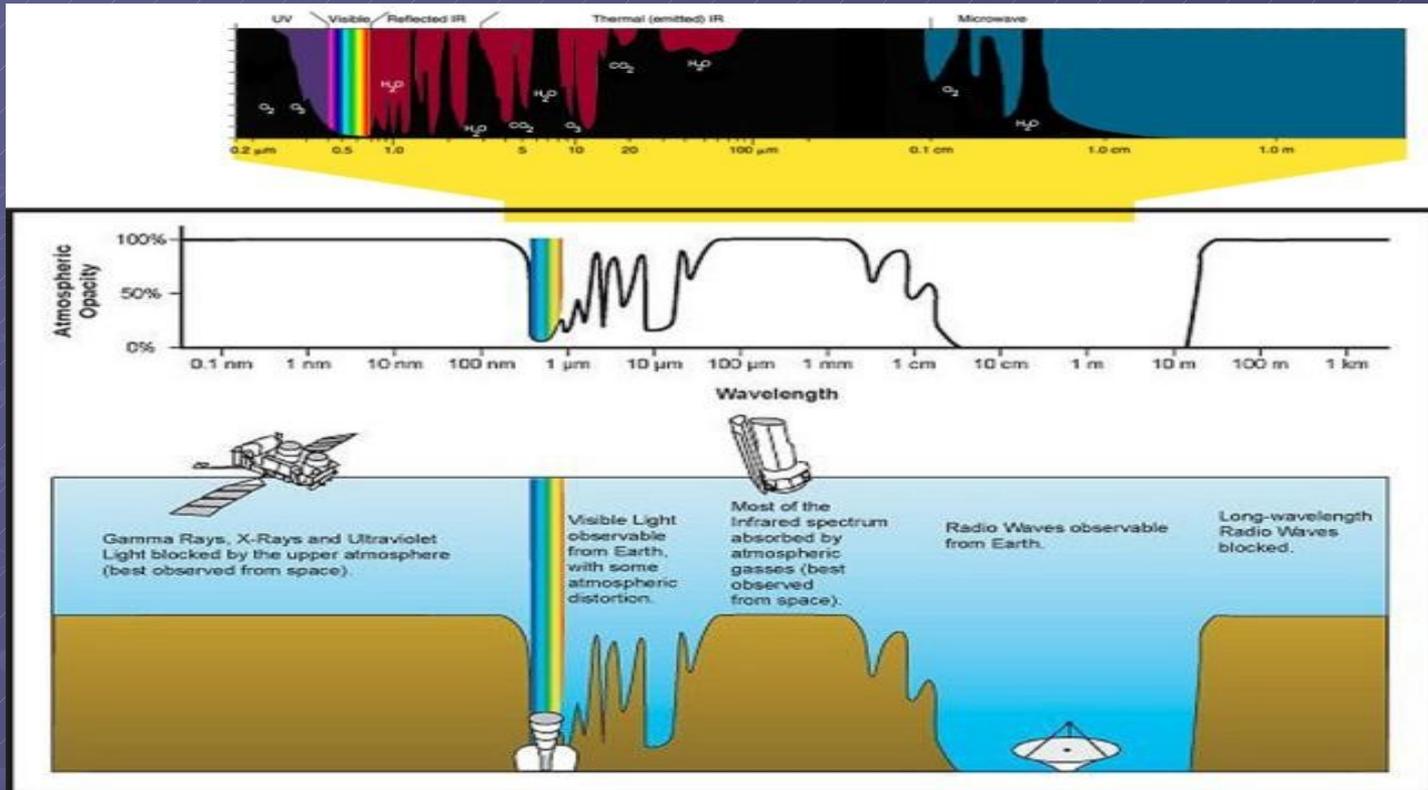
# Radio Techniques

## Single Dish Radioastronomy

### Part I

Dipanjan Mitra  
NCRA

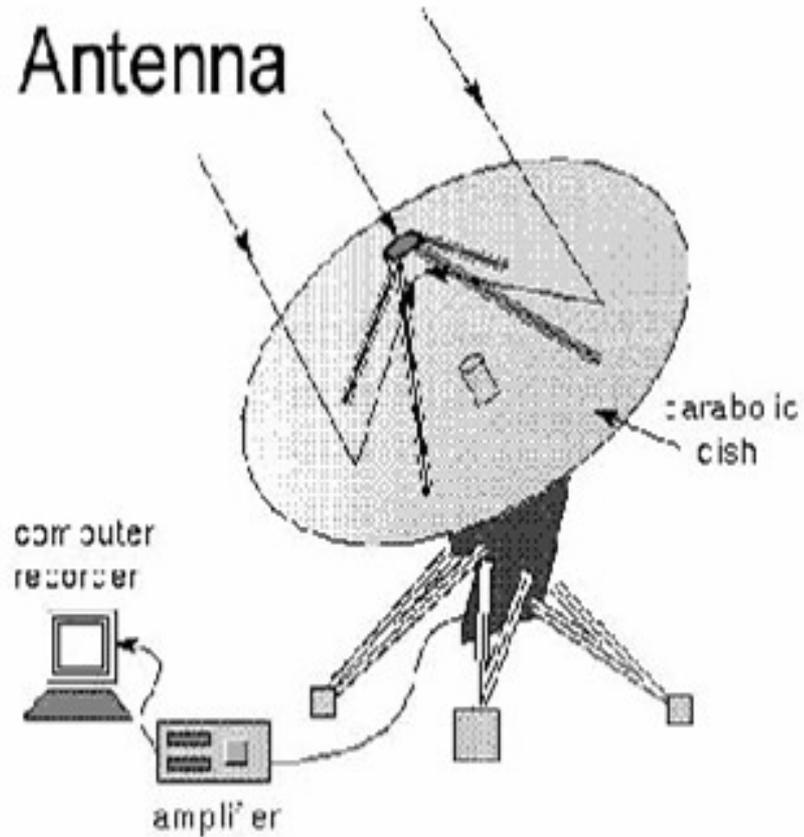
# ATMOSPHERIC WINDOW



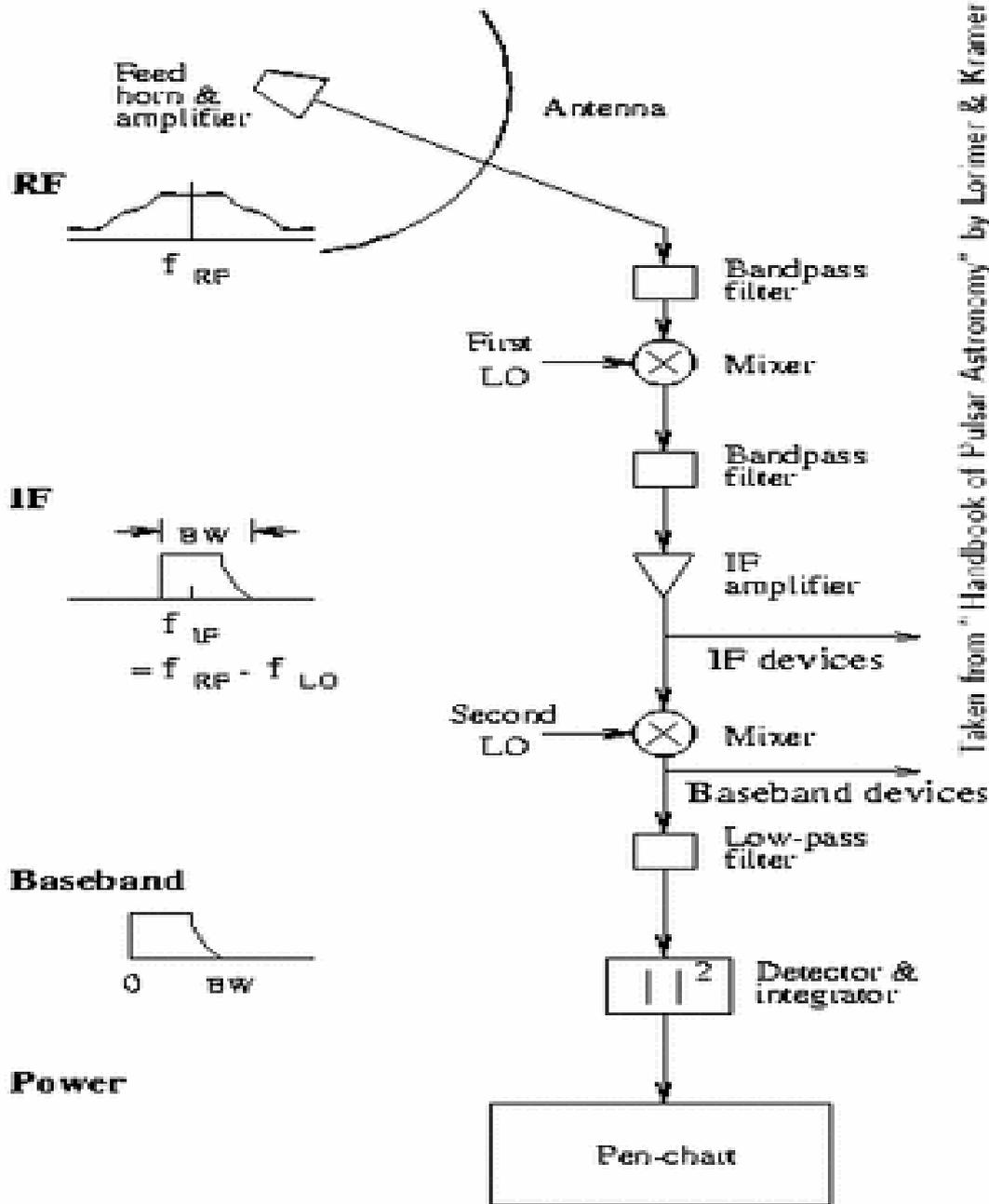
HIGH FREQUENCY CUTOFF IN RADIO IS: 600 GHz (or 0.5 mm)

LOW FREQUENCY CUTOFF IS : 15 MHz (or 20 m)

# Schematic of a Single Dish Radio Telescope



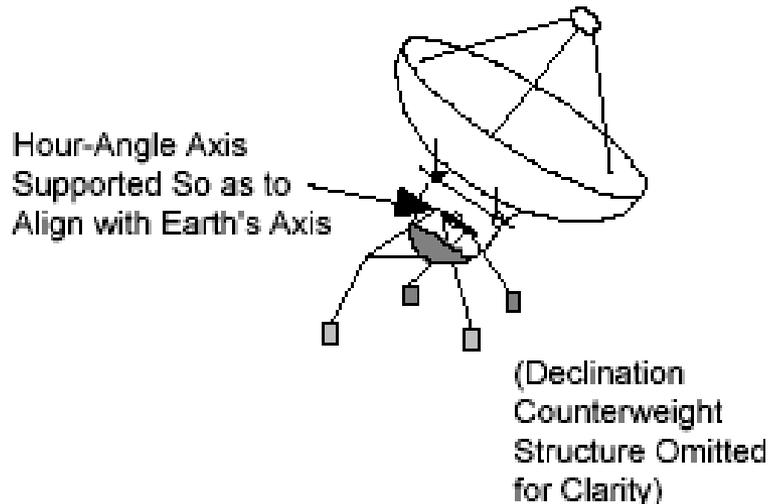
A radio telescope reflects radio waves to a focus at the antenna.



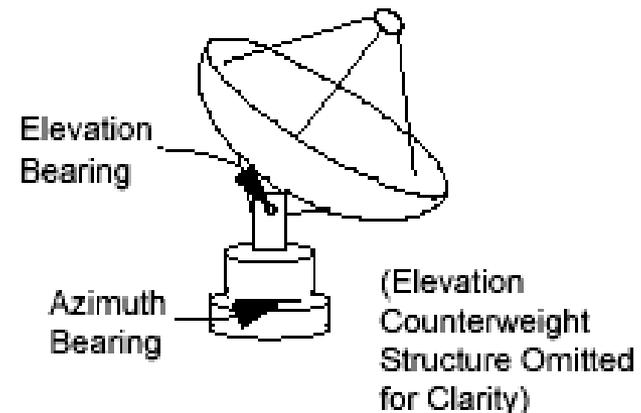
Taken from "Handbook of Pulsar Astronomy" by Lorimer & Kramer

# Mounting of Antennas

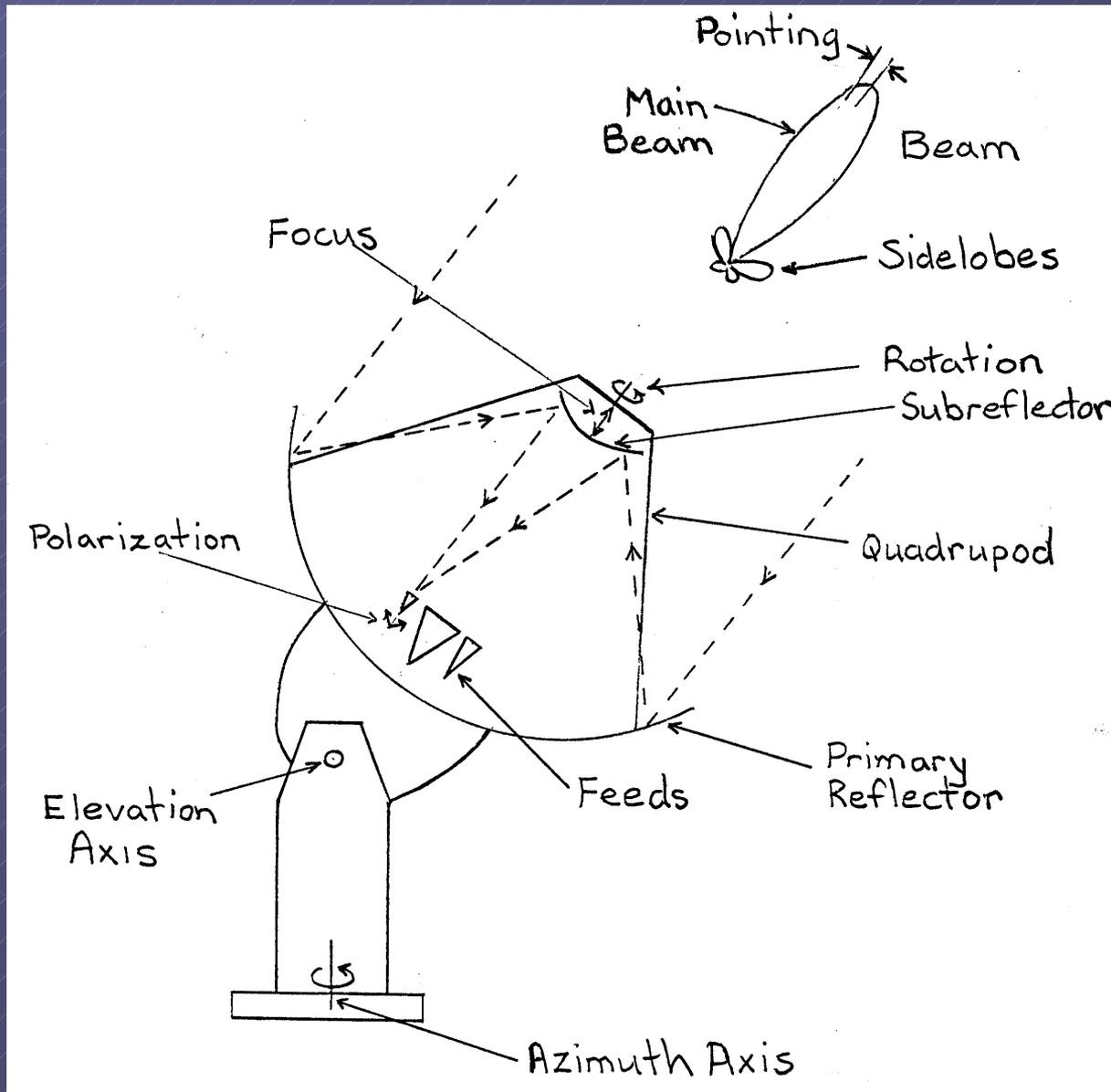
**HA-DEC Mount**



**AZ-EL Mount**



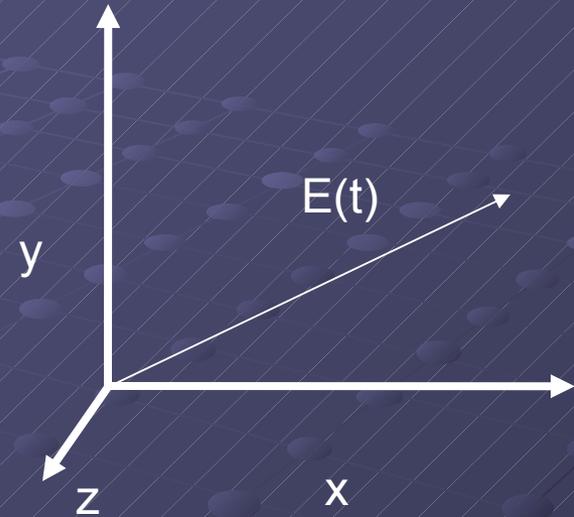
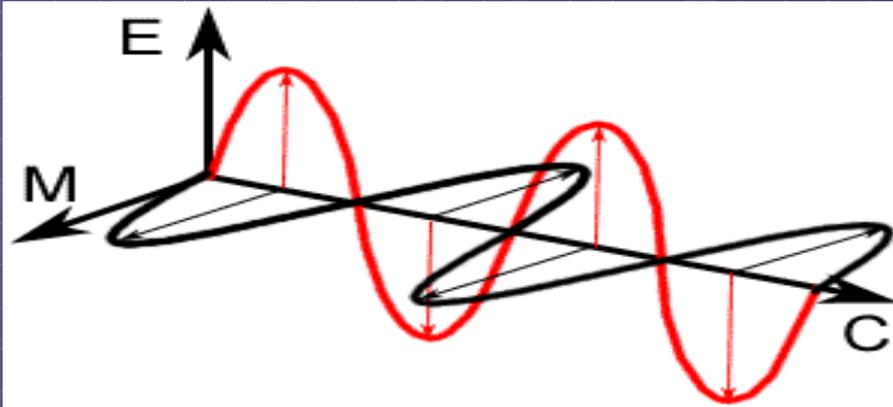
# Key elements of an Antenna



# Nature of the radio emission of a Celestial source

- At radio frequencies the EM radiation can be treated as waves
- EM waves are generated by acceleration of charged particles
- Inverse square law: Flux of the source decrease as  $1/R^2$
- Celestial sources are so far away that we can ignore the curvature of the wave front (plane wave)

# Plane Wave Basics



$$E_x = A_x \cos(2\pi\nu t + \delta_x)$$

$$E_y = A_y \cos(2\pi\nu t + \delta_y)$$

- For  $A_x$  or  $A_y = 0$ , linearly polarized light.
- For  $A_x = A_y$ ,  $\delta_x = 0$ ,  $\delta_y = -\pi/2$ , circularly polarized light

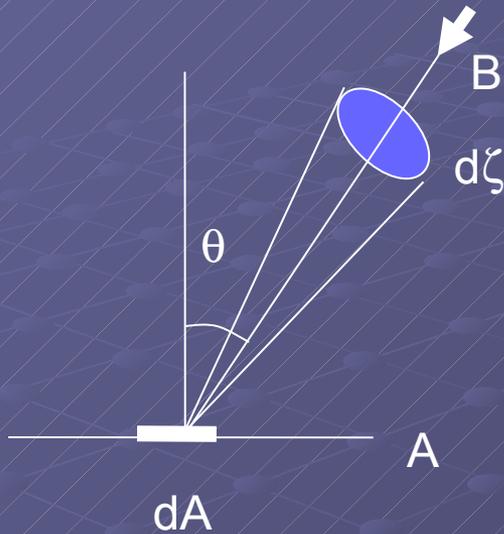
# Plane wave cont...

Variation of the Plane wave in space and time can be described as

$$E(z,t) = A \cos (2\pi\nu t - k z)$$

$$\text{or } E = \text{Real} ( A e^{j(2\pi\nu t - kz)} )$$

# Some Basic Definitions in Radioastronomy



infinitesimal power from a solid angle  $d\zeta$  is  
 $dW = B \cos(\theta) d\zeta dA dv$  watts

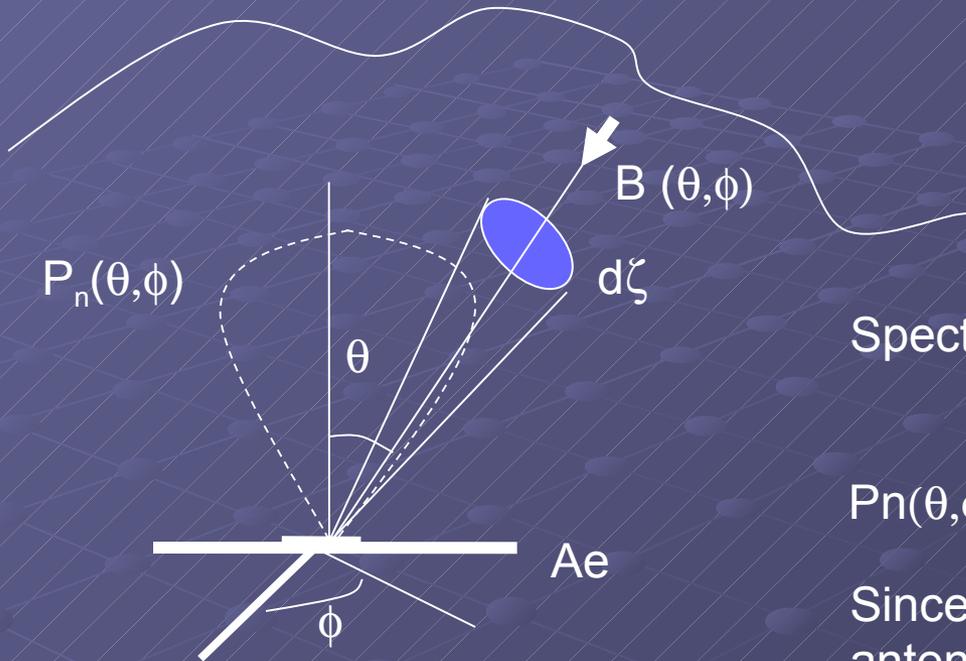
$B =$  Brightness , watt  $m^{-2} Hz^{-1} rad^{-2}$

$$\text{Power } W = A \int_{\nu} \int_{\zeta} B \cos(\theta) d\zeta dv$$

Spectral power :  $dw = B \cos(\theta) d\zeta dA$  watt  $Hz^{-1}$

For a constant brightness the spectral power  $w = \pi A B$

# Definitions continued . . . . .



$$\text{Spectral Power : } w = Ae \int B(\theta, \phi) P_n(\theta, \phi) d\zeta$$

$P_n(\theta, \phi)$  is the power pattern of the antenna

Since the radiation received by the antenna is of incoherent type and since any antenna is only responsive to 1 polarization the spectral power is

$$w = \frac{1}{2} Ae \int B(\theta, \phi) P_n(\theta, \phi) d\zeta$$

# Flux Density of a source

$$S = \int B(\theta, \phi) P(\theta, \phi) d\zeta$$

The unit of flux is jansky (Jy) which is equal to  $10^{-26}$  watt  $m^{-2}$   $Hz^{-1}$

Case 1) For a point source, where the source is much smaller than the antenna beam

$$S = \int B(\theta, \phi) d\zeta$$

Case 2) When the source is much larger than the main lobe of the antenna and the brightness is constant over the main lobe

$$S = B(\theta, \phi) \int P(\theta, \phi) dz \sim B(\theta, \phi) \zeta_m$$

# Concept of Temperature

Radiation from a blackbody is described by Planks law

$$B(\nu) = \left( \frac{2h\nu^3}{c^2} \right) \frac{1}{(e^{h\nu/kT} - 1)} \text{ Watt m}^{-2} \text{ Hz}^{-1} \text{ rad}^{-2}$$

For a typical radio frequency like  $10^9$  Hz,  $h\nu/k \sim 0.048$ , Hence

$$e^{h\nu/kT} \sim 1 + h\nu / kT \text{ or, } B(\nu) = \left( \frac{2\nu^2}{c^2} \right) k T = 2 k T / \lambda^2 \text{ or } T = \left( \frac{\lambda^2}{2 k} \right) B(\nu)$$

This approximation is called the Rayleigh-Jeans Approximation of the Plank spectrum.

The brightness temperature of a source is hence defined as  $T_b = \left( \frac{\lambda^2}{2 k} \right) B(\nu)$

In general  $T_b$  has no relation to the physical temperature of the source (except some sources in the sky which are emitting as a black body)

Why do we think that pulsar radio emission is coherent?

Pulsars have high brightness temperature

$$T_b = (c^2 / 2\nu^2 k) I_\nu$$

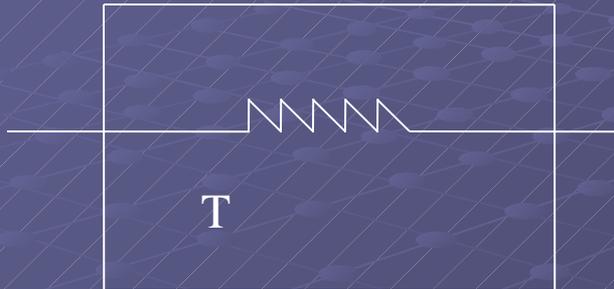
$$F_\nu = I_\nu \Omega, \quad \Omega \sim l^2/d^2$$

$$T_b \sim (3.1 \times 10^{23}) \nu_9^{-2} F_\nu (\text{mJy}) d^2 (\text{kpc}) I_6^{-2}$$

$$T_b \sim 10^{25} \text{ --- } 10^{30} \text{ k}$$

this is extremely high !!!

# Nyquist Theorem and Noise Temperature



A resistor put in a thermal bath of temperature  $T$  will have an output power per unit frequency given by

$$P = k T , \text{ called the Nyquist formula}$$

When radiation falls in an antenna, the power absorbed is expressed in temperature units. The power available in the antenna due to a source the sky is termed as the antenna temperature  $T_a$

$$T_a = P_a / k$$

The power introduced in a radio telescope due to several noise contributions (e.g from sky, ground, receiver etc) is called the system temperature  $T_{\text{sys}}$

$$T_{\text{sys}} = P_{\text{sys}} / k$$

# Connecting Flux and Temperature

Recall that the spectral power is

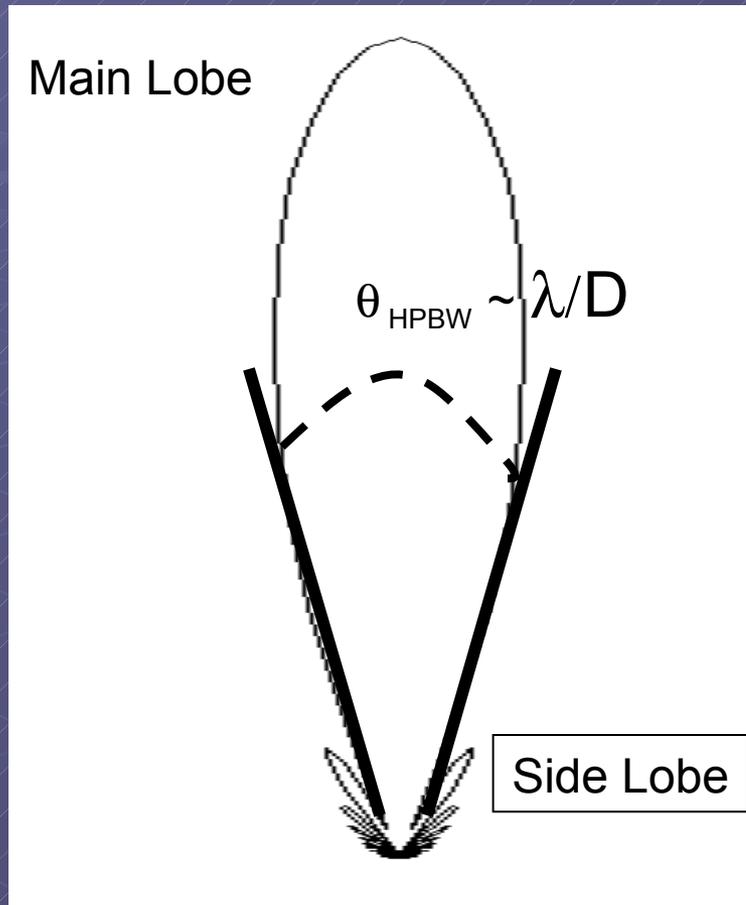
$$w = \frac{1}{2} A_e \int B(\theta, \phi) P_n(\theta, \phi) d\zeta$$

again  $w = k T$  as per Nyquist theorem

Hence we get a relation between temperature and flux as

$$kT = \frac{1}{2} A_e S$$

# Beam of an Antenna



## Reciprocity

The pattern of an antenna is same whether it is used as a transmitting or receiving antenna, i.e if the antenna emits efficiently in one direction it will receive efficiently in that direction.

# Antenna Patterns

The effective aperture of the antenna

$$A_e(\theta, \phi) = \frac{\text{Power density available at the antenna}}{\text{flux density of the wave incident on the antenna}}$$

The power pattern of the antenna

$$P(\theta, \phi) = \frac{A_e(\theta, \phi)}{A_e^{\max}}$$

$$\text{Directivity of the antenna, } D(\theta, \phi) = \frac{4\pi P(\theta, \phi)}{\int P(\theta, \phi) d\zeta} = \frac{\text{Power emitted in } (\theta, \phi)}{\text{Total power emit./ } 4\pi}$$

Gain = Power emitted in  $(\theta, \phi)$  / Total power input /  $4\pi$

cont.....

.....  
 Recall the expression for spectral power:  $w = \frac{1}{2} A_e \int B(\theta, \phi) P_n(\theta, \phi) d\zeta$

$$T_a(\theta', \phi') = A_e^{\max} / \lambda^2 \int T_b(\theta, \phi) P(\theta - \theta', \phi - \phi') \sin(\theta) d\theta d\phi$$

Consider an antenna terminated in a resistor:

In thermal equilibrium

The power  $P_{-}(R \rightarrow A) = kT$

$$P_{-}(A \rightarrow R) = (A_e^{\max} kT / \lambda^2) \int P(\theta, \phi) d\zeta \quad \text{or} \quad A_e^{\max} = \lambda^2 / \int P(\theta, \phi) d\zeta$$

$$A_e = A_e^{\max} P(\theta, \phi) = \frac{\lambda^2 P(\theta, \phi)}{\int P(\theta, \phi) d\zeta} \quad \text{or} \quad D(\theta, \phi) = (4\pi / \lambda^2) A_e(\theta, \phi)$$

Gain of the antenna is connected to the directivity through a constant factor.

$$\varepsilon = Ae^{\max} / Ag$$

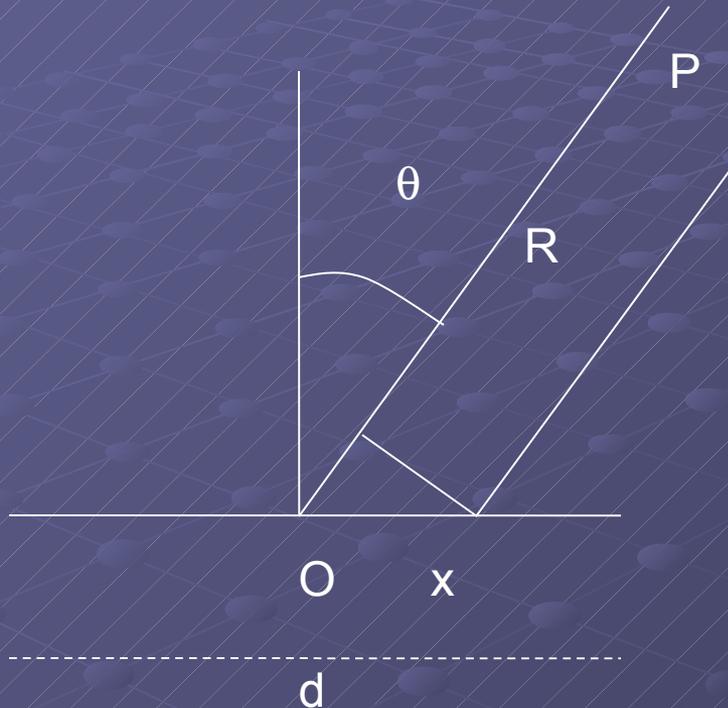
## F/D ratio for an antenna

In a antenna feed system the feed should ideally illuminate the antenna with a uniform beam that only illuminates the reflecting surface

The angle subtended by the feed as seen by the reflector is given by the Focus / Diameter , F/D ratio.

When  $F/D \sim 0.38$ , the edge of the dish is about 64 deg as seen by the feed and the efficiency is close to 100 %

# Computing Antenna Patterns



The beam is the fourier transform the of the aperture

The total electric field at  $P$  due to a electric field distribution  $e(x)$  is

$$E(R, \theta) = \int_{-d/2}^{d/2} \frac{e(x)}{R^2} e^{-jk_{\mu}x} dx$$

The un-normalised power has the form of a fourier transform

$$F(\mu) = \int_{-inf}^{+inf} e1(x) e^{-jk_{\mu}x} dx$$

# Rules of Fourier Transform

Linearity, if  $G_1(\mu) = F[g_1(x)]$  and  $G_2(\mu) = F[g_2(x)]$   
then:  $G_1(\mu) + G_2(\mu) = F[g_1(x) + g_2(x)]$

Inverse: if  $G(\mu) = \int_{-\infty}^{\infty} g(x) e^{-jk\mu x} dx$ ,

then

$$g(x) = \int_{-\infty}^{\infty} G(\mu) e^{jk\mu x} dx$$

Phase Shift:  $G(\mu - \mu_0) = F [ g(x) e^{-j2\pi\mu_0 x} ]$ , This is the principle used to steer the antenna beam

# Parsevals Theorem

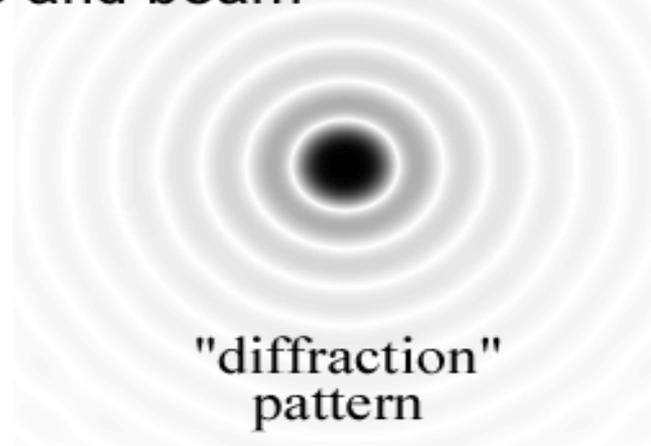
$$\int_{-\infty}^{\infty} |G(\mu)|^2 d\mu = \int_{-\infty}^{\infty} |g(x)|^2 dx$$

Area:  $G(0) = \int_{-\infty}^{\infty} g(x) dx$

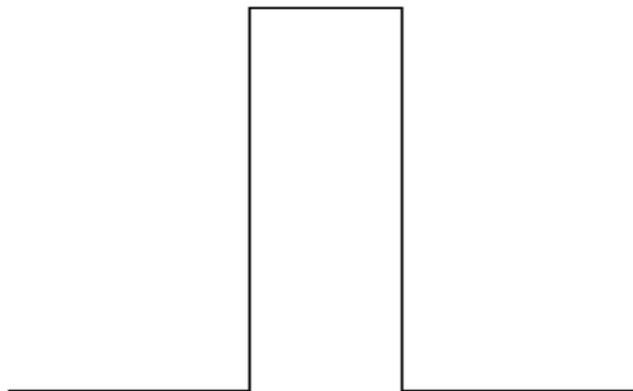
# Fourier Transform relationship of aperture and beam



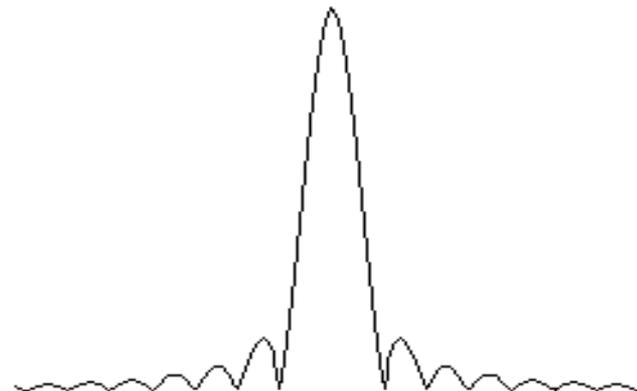
circular  
aperture



"diffraction"  
pattern



cross-section  
of aperture



cross-section  
of diffraction pattern

# Factors affecting Antenna Performance

- Reflector surface efficiency
- Blockage by the feed
- Spillover
- Surface accuracy (Ruze loss)
- Feed illumination efficiency (F/D ratio)
- Low sidelobe pattern
- Pointing

# Reber's 31.4 ft parabolic reflector





# Greenbank 100-m telescope





300-m telescope in Arecibo, Puerto Rico

# GMRT 45 metre dish



Note the light , see-through structure reducing wind forces, weight, and cost

Ooty Radio Telescope 530mx30m parabolic cylinder operating at 325MHz



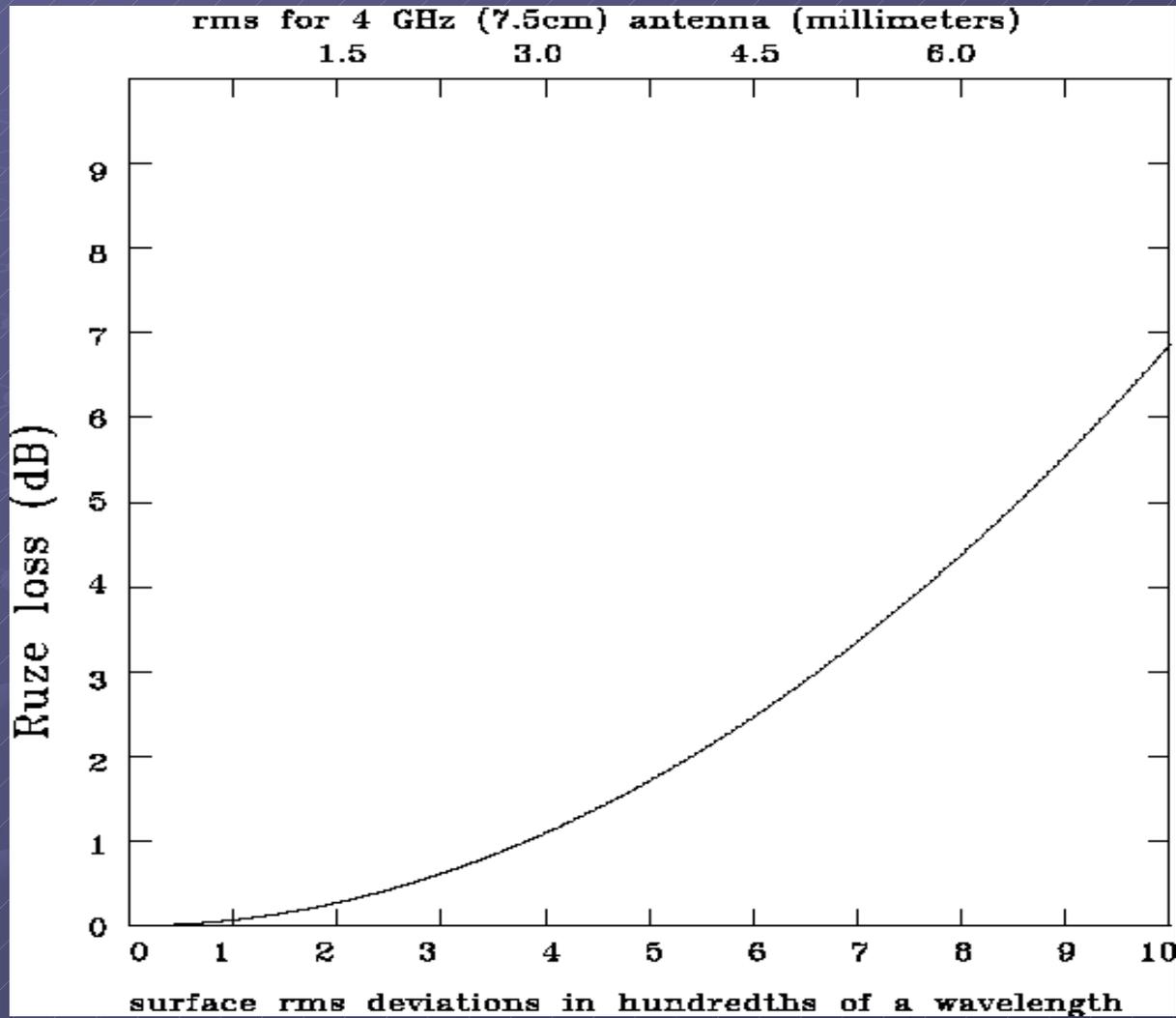
# Gauribidnur Array at 30 MHz



$$L = \exp \left( - \left( 4 \pi \sigma / \lambda \right)^2 \right)$$

$$\sigma = \lambda / 16$$

causes a  
loss of 0.5



# Antenna (Collector)

- Area  $\sim 10^4 \text{ m}^2$  ,  $P = \frac{1}{2} A S \Delta\nu \sim 10^{-15} \text{ watts}$
- Radio Telescopes are diffraction limited,  
Resolution  $\sim \lambda / D \sim 0.5 \text{ deg}$  (for wave 1 m and  $D=100$ )  
Human eye  $\sim 20''$  , Ground based optical tel:  $1''$   
Directivity  $\sim \text{Gain}(\Theta) / \epsilon$
- Surface accuracy  $\sim \lambda/10$  , such that signal can add coherently at the focus.  
Mesh surface possible at low freq, like in GMRT, and is cost effective

# Feed System

- At the feed the EM wave is converted into electrical signal in the cable
- Feed are resonant devices like dipoles or horns with  $\delta\nu/\nu = 10 - 20 \%$
- Feed is small and hence have large angular beam. So to avoid picking up ground radiation feed is designed to have  $1/10^{\text{th}}$  gain at the edge of the dish

This leads to effective loss of collecting area

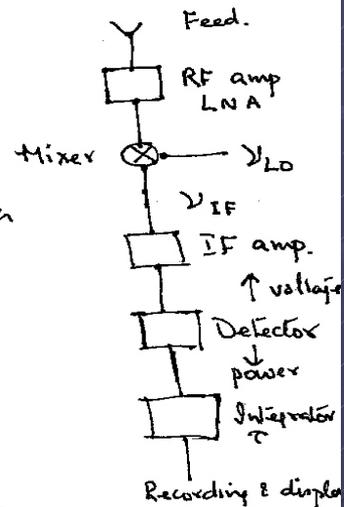
- In GMRT gain = 0.3 K / Jy, Ground radiation is 300 K, efficiency is around 70 %

## Electronics

- Super heterodyne receivers - similar to household radio
- Radio source signals are very weak - amplification of  $10^{6-8}$
- Amplifier characterised by gain  $G$ .
- Due to collision of electrons in the amplifier, it ~~corrupts~~ <sup>adds noise to</sup> the input signal before amplification

$$\text{Output power of amplifier} = G \times k (T_{\text{ant}} + T_{\text{Receiver}})$$

- $T_{\text{Receiver}} \sim 20 \text{ K to } 100 \text{ K}$
- Low Noise RF amplifiers made of using High Electron Mobility Transistors and other special devices
- Often (at cm & mm  $\lambda$ ) cooled to Liquid Nitrogen w lower temperatures
- At low frequencies ( $< 300 \text{ MHz}$ ) LNAs not critical - galactic background emission ( $T_{\text{Back Sky}} \sim 40 \text{ K @ } 325 \text{ MHz}$   $\propto \nu^{-2.5}$ ) is the dominant source of noise
- RF amplifier determines S/N - other amplifiers in chain don't matter



### Heterodyning - Mixer

For QM signal  $A(t) e^{i[\omega t + \phi(t)]}$  all the information is in  $A(t), \phi(t)$  - nothing in  $\omega$

If we multiply the QM signal by a pure sinusoid at  $\nu_{LO}$

$$\text{Output} = A(t) e^{i[(\omega - \omega_{LO})t + \phi(t)]} = A(t) e^{i[\omega_{IF}t + \phi(t)]}$$

Mixer is a non linear device that achieves this multiplication - [avoid cable loss - standard components]

## Distribution of Input voltage

$$P(V) = \frac{1}{(2\pi)^{1/2}\sigma} \exp\left(\frac{-V^2}{2\sigma^2}\right),$$

After detection  $V_o = V^2$ , assuming  $\sigma = 1$

$$P_o(V_o)dV_o = 2P(V)dV$$

for  $V \geq 0$ . Since  $dV_o = 2VdV$ ,  $dV/dV_o = V_o^{-1/2}/2$  and

$$P_o(V_o) = \frac{1}{(2\pi)^{1/2}} V_o^{-1/2} \exp(-V_o/2)$$

$$\langle V_o \rangle = \sigma^2.$$

RMS =  $2^{1/2}$  times the mean output power